

**SEARCH FOR SOLAR AXIONS IN THE CAST  
EXPERIMENT USING THE MICROME GAS  
DETECTOR**

PhD THESIS

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# Outline

## I. Theory

The strong CP problem, the dynamic Peccei–Quinn solution, axion and its properties, axion–photon oscillation.

## II. CAST

Solar axions, principle of the experiment, the magnet and the detectors.

## III. Micromegas

General description and properties, the micromegas of CAST, typical performance.

## IV. Analysis & Results

Simulation, 2003 data, 2004 data.



# THEORY



# The Strong CP Problem

$$\vec{G}_{\mu\nu}^{\tilde{z}} = \frac{1}{2} \epsilon_{\mu\nu\alpha\beta} \vec{G}^{\alpha\beta}$$

$$L_{\bar{\theta}} = \frac{g_s^2 \bar{\theta}}{32\pi^2} \vec{G}^{\mu\nu} \cdot \vec{G}_{\mu\nu}^{\tilde{z}}$$

$$\vec{G}^{\mu\nu} \cdot \vec{G}_{\mu\nu}^{\tilde{z}} \quad \partial^\mu K_\mu$$

$$K_\mu = 2\epsilon_{\mu\alpha\beta\gamma} \vec{A}^a \cdot \left\{ \partial^\beta \vec{A}^\gamma + \frac{g_s}{3} (\vec{A}^\beta \times \vec{A}^\gamma) \right\}$$

$$|\theta\rangle = \sum_n e^{in\theta} |n\rangle$$

$$\bar{\theta} = \theta + \text{Arg}(\det(M_q))$$

Neutron electric dipole moment

$$\left. \begin{array}{l} d_n^{\text{theory}} \approx \bar{\theta} \times 10^{-15} \text{ e} \cdot \text{cm} \\ d_n^{\text{exp}} < 0.63 \times 10^{-25} \text{ e} \cdot \text{cm} \end{array} \right\} \Rightarrow \bar{\theta} \sim 10^{-10}$$

Extra term in the QCD Lagrangian:

- α. CP violating.
- β. Total divergence (Bardeen current).
- γ. Not affecting the perturbative behavior of the theory.
- δ. Non zero contribution (vacuum topology).

Why so small;

The Standard Model **DOES NOT** provide an answer.



# The Peccei-Quinn mechanism

**PQ Symmetry:** global, axial, beyond the Standard Model.

**Spontaneous symmetry breaking at scale  $f_a$**

**Axion:** Goldstone boson, pseudoscalar, neutral.

$$\mathcal{L}_a = C \frac{g_s^2}{32\pi^2 f_a} \alpha(\mathbf{x}) \vec{G}^{\mu\nu} \cdot \vec{G}_{\mu\nu}$$

$$V_{\text{eff}}(\alpha) = \frac{1}{2} K \left( \bar{\theta} + C \frac{\alpha(\mathbf{x})}{f_a} \right)^2$$

$$\langle \alpha \rangle = -\frac{f_a \bar{\theta}}{C} \quad m_a^2 = \left( \frac{\partial^2 V_{\text{eff}}}{\partial^2 \alpha} \right)_{\langle \alpha \rangle}$$

The vacuum expectation value of the axion field eliminates the QCD term at the cost of axion appearance.



# Axion models

## PQWW

- $f_\alpha = f_{ew}$
- Two Higgs fields.
- Fermions: PQ charge.

## DFSZ

- $f_\alpha \gg f_{ew}$
- Two Higgs fields and one scalar field.
- Fermions: PQ charge.

## KSVZ

- $f_\alpha \gg f_{ew}$
- One Higgs field, one scalar field, one exotic quark with PQ charge.

Experimentally excluded.

$$\Phi'_u = e^{i\alpha\Gamma_u} \Phi_u, \quad \Phi'_d = e^{i\alpha\Gamma_d} \Phi_d$$

$$u'_L = e^{\frac{1}{2}i\alpha\Gamma_u} u_L, \quad u'_R = e^{-\frac{1}{2}i\alpha\Gamma_u} u_R$$

$$d'_L = e^{\frac{1}{2}i\alpha\Gamma_d} d_L, \quad d'_R = e^{-\frac{1}{2}i\alpha\Gamma_d} d_R$$

$$\Phi'_u = e^{i\alpha X_u} \Phi_u, \quad \Phi'_d = e^{i\alpha X_d} \Phi_d$$

$$\sigma' = e^{i\alpha X_\sigma} \sigma, \quad X_u + X_d = -2X_\sigma = 1$$

$$u'_L = e^{\frac{1}{2}i\alpha X_u} u_L, \quad u'_R = e^{-\frac{1}{2}i\alpha X_u} u_R$$

$$d'_L = e^{\frac{1}{2}i\alpha X_d} d_L, \quad d'_R = e^{-\frac{1}{2}i\alpha X_d} d_R$$

$$\psi'_L = e^{\frac{1}{2}i\alpha} \psi_L, \quad \psi'_R = e^{-\frac{1}{2}i\alpha} \psi_R$$

$$\sigma' = e^{i\alpha} \sigma$$

Hadronic axion



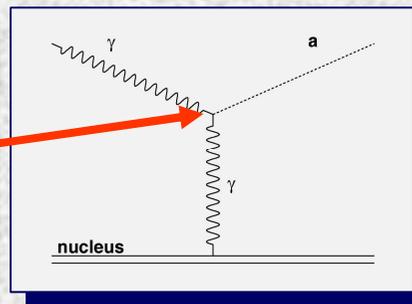
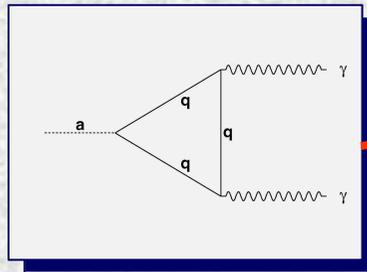
# Properties

**Electromagnetic interaction**

$$L_a = \frac{1}{4} g_{a\gamma\gamma} \alpha F^{\mu\nu} \tilde{F}_{\mu\nu}$$

$$g_{a\gamma\gamma} = C \frac{e^2}{2\pi^2 f_a}$$

$$C = \sum_{\text{fermions}} Q_{\text{em}}^2 Q_{\text{PQ}}$$



**Primakoff effect**

$$\frac{d\sigma_{a \leftrightarrow \gamma}}{d\Omega} = \frac{g_{a\gamma\gamma}^2 Z^2 e^2}{32\pi^2} \frac{|\vec{p}_a \times \vec{p}_\gamma|^2}{|\vec{p}_a - \vec{p}_\gamma|^4}$$

**Axion lifetime**

$$\tau(a \rightarrow \gamma\gamma) \sim \frac{10^{23}}{[m_a (\text{eV})]^5} \text{sec}$$

**Axion mass**

$$m_a = N_g m_\pi \cdot \frac{f_\pi}{f_a} \cdot \frac{\sqrt{m_u m_d}}{m_u + m_d}$$

**Axion-Photon conversion  
in the Coulomb field of a  
nucleus.**

**$f_a \gg 1\text{GeV}$**

**Invisible Axion**

**Number of quark  
generations carrying  
PQ charge.**



# Axion – Photon oscillation

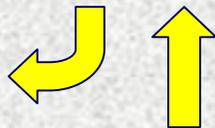
## Assumptions

- vacuum,
- weak interaction:  $g_{\alpha\gamma\gamma} \ll \frac{m_a^2}{BE}$
- transverse magnetic field:  $\vec{p}_a \perp \vec{B}$
- relativistic limit:  $E \gg m_a$

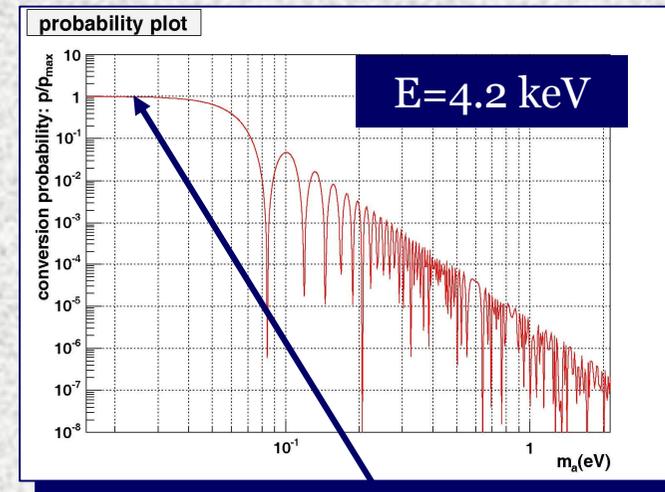
$$L = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} (\partial_\mu \alpha \partial^\mu \alpha - m_a^2 \alpha^2) + \frac{1}{4} g_{\alpha\gamma\gamma} \alpha F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$(\nabla^2 + m_a^2) \alpha - g_{\alpha\gamma\gamma} \frac{\partial \vec{A}}{\partial t} \cdot \vec{B} = 0$$

$$\nabla^2 \vec{A} + g_{\alpha\gamma\gamma} \frac{\partial \alpha}{\partial t} \vec{B} = \vec{0}$$



$$F_{\mu\nu} = F_{\mu\nu}^{\text{ext}} + \partial_\mu A_\nu - \partial_\nu A_\mu$$



$$P_{\alpha \rightarrow \gamma}(z) = \frac{B^2 g_{\alpha\gamma\gamma}^2}{q^2} \sin^2\left(\frac{qz}{2}\right) \xrightarrow{q \rightarrow 0} \frac{1}{4} B^2 g_{\alpha\gamma\gamma}^2 z^2$$

$$q = \frac{m_a^2}{2E} \text{ (vacuum momentum transfer)}$$



# CAST

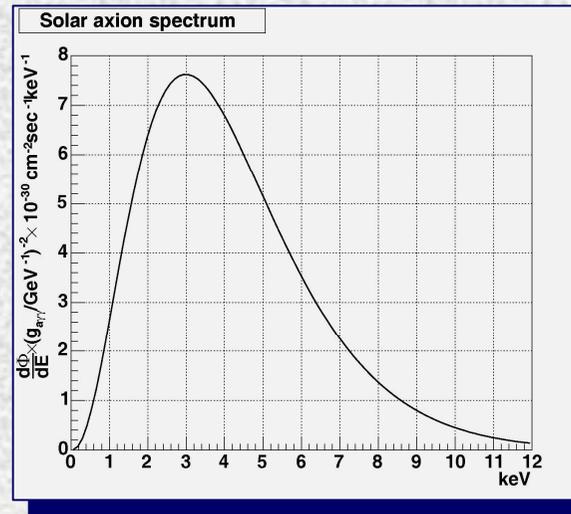
## Cern Axion Solar Telescope



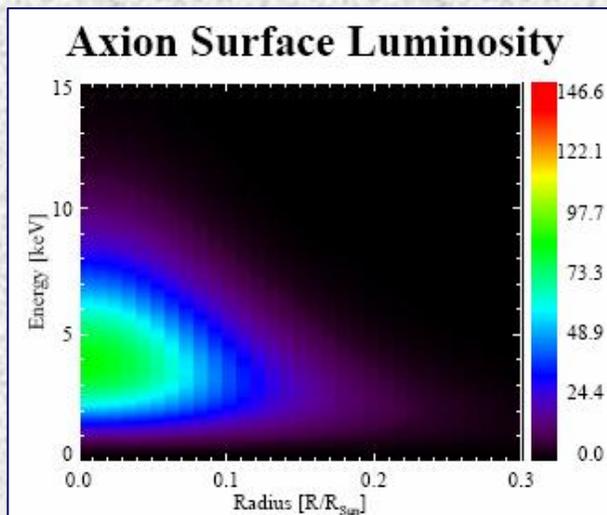
# Solar Axions

## Assumptions

- Hadronic Axions.
- Primakoff Effect.
- Charge screening.
- Standard Solar Model.



Axion spectrum on Earth



$$\frac{d\Phi_{\alpha}}{dE} = 6.02 \times 10^{10} \cdot g_{10}^2 \cdot E^{2.481} \cdot e^{-E/1.205} \text{ cm}^{-2} \text{ s}^{-1} \text{ keV}^{-1}$$

$$L_{\alpha} = 1.85 \times 10^{-3} \cdot g_{10}^2 \cdot L_{\odot}$$

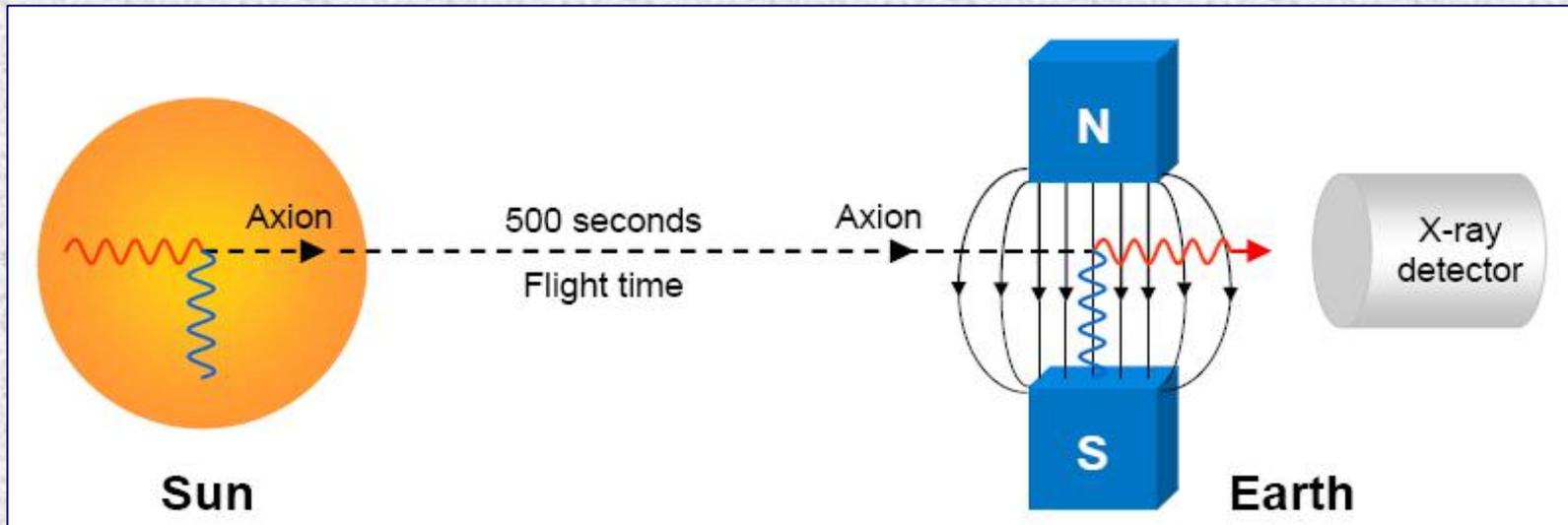
$$\langle E_{\alpha} \rangle = 4.2 \text{ keV}$$

$$\Phi_{\alpha} = 3.67 \times 10^{11} \cdot g_{10}^2 \cdot \text{cm}^{-2} \text{ s}^{-1}$$

$$g_{10} = g_{\alpha\gamma\gamma} \times (10^{10} \text{ GeV})$$



# Detection Principle



Axions are produced in the Solar core through the Primakoff effect and travel freely to the Earth where they are converted to photons, interacting with a transverse magnetic field which can align to the Sun (Helioscope). If there is an excess of photons when tracking the Sun, it is an indication of axion existence.



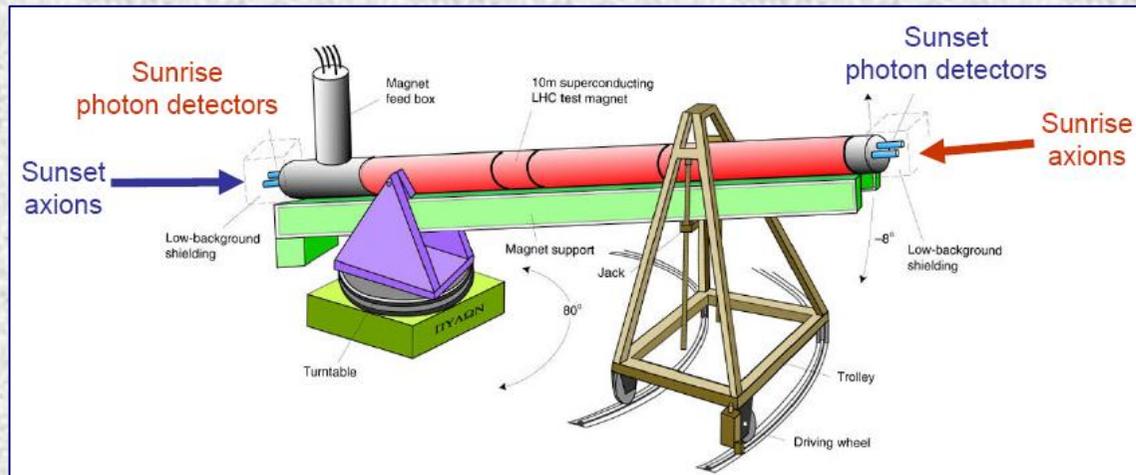
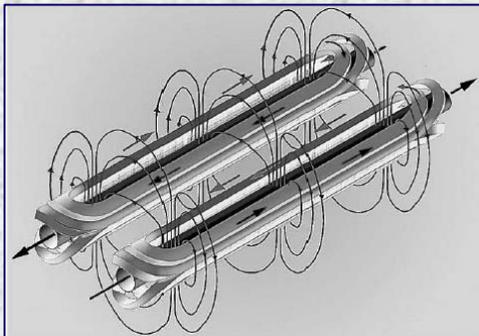
# Experimental apparatus



- Prototype, dipole LHC magnet.
- Superconducting:  
 $T=1.8\text{K}$ ,  $I=13,000\text{A}$ .
- $B=9\text{T}$ ,  $L=9.26\text{m}$ .
- $\pm 8^\circ$  vertical motion,  $80^\circ$  horizontal motion.
- 3h Sun tracking daily.

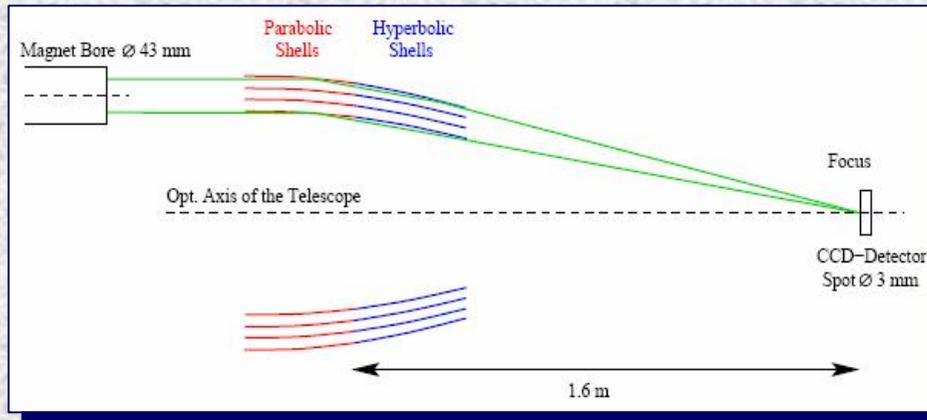
**Phase I:** pipes in vacuum.

**Phase II:** pipes filed with gas ( $\text{He}_3$ ,  $\text{He}_4$ ).





# Detectors



## Telescope and CCD

- Wolter I optical device, prototype of ABRIXAS space mission. 27 coaxial gold-plated shells from Ni. Axionic image of the Sun.
- Charged Coupled Device (silicon detector). Pixel size:  $150 \times 150 \mu\text{m}^2$ .
- Spot: 3mm diameter, huge increase of CAST sensitivity.
- Covers one magnet bore and tracks the Sun during sunrise.

## TPC

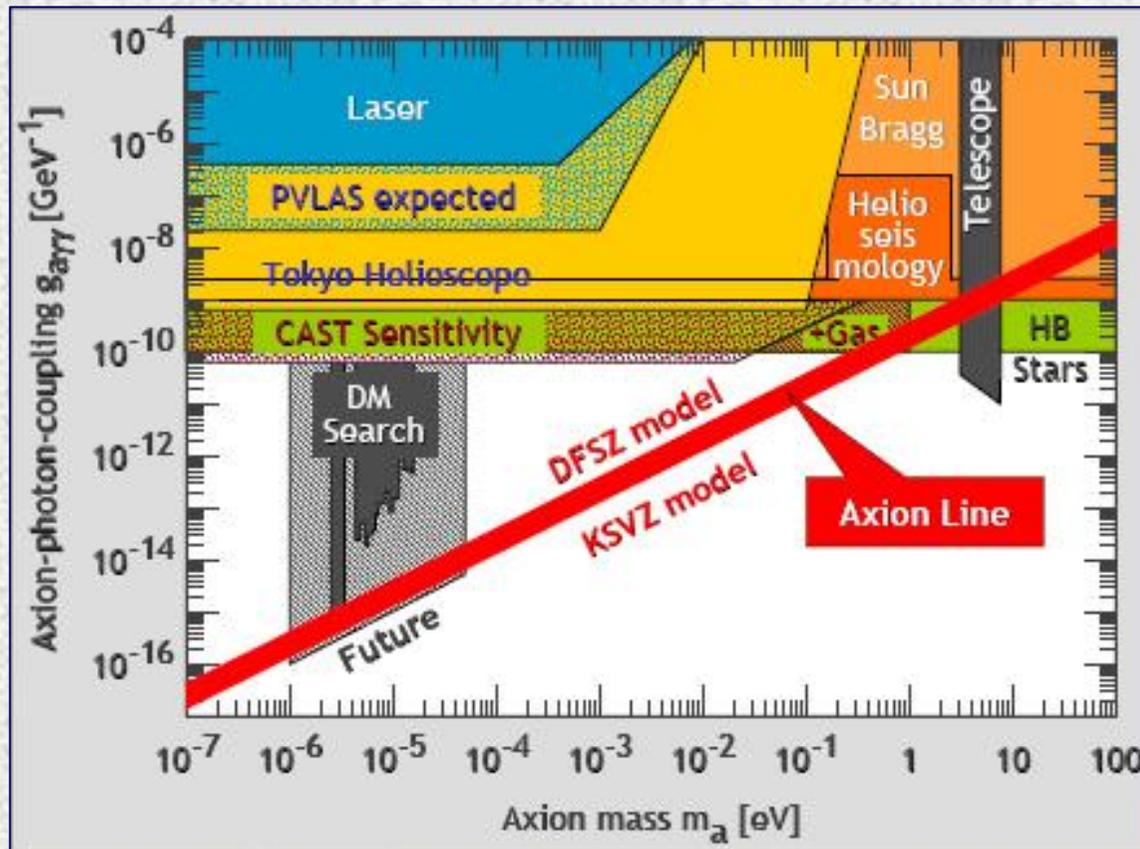
- Time Projection Chamber: conventional gas detector with 48 anode wires and 96 cathode wires, 3mm apart.
- Shielded. Constructed by Zaragoza University.
- Covers two magnet bores and tracks the Sun during sunset.

## MICROME GAS

- New technology, parallel plate gas detector with microstructure and spatial sensitivity.
- Constructed by Saclay and NCSR "Demokritos".
- Covers one magnet bore and tracks the Sun during sunrise.



# Axionic parametric space



## Sources

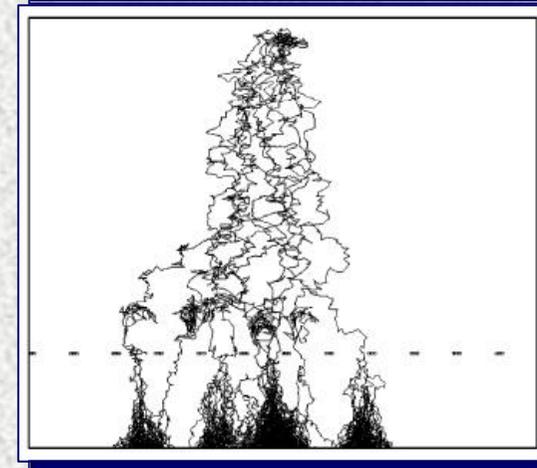
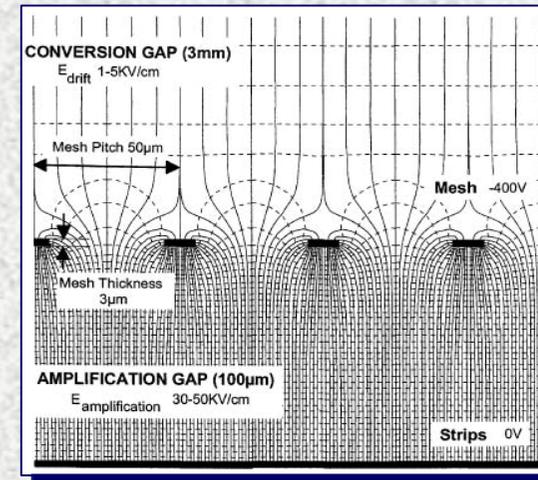
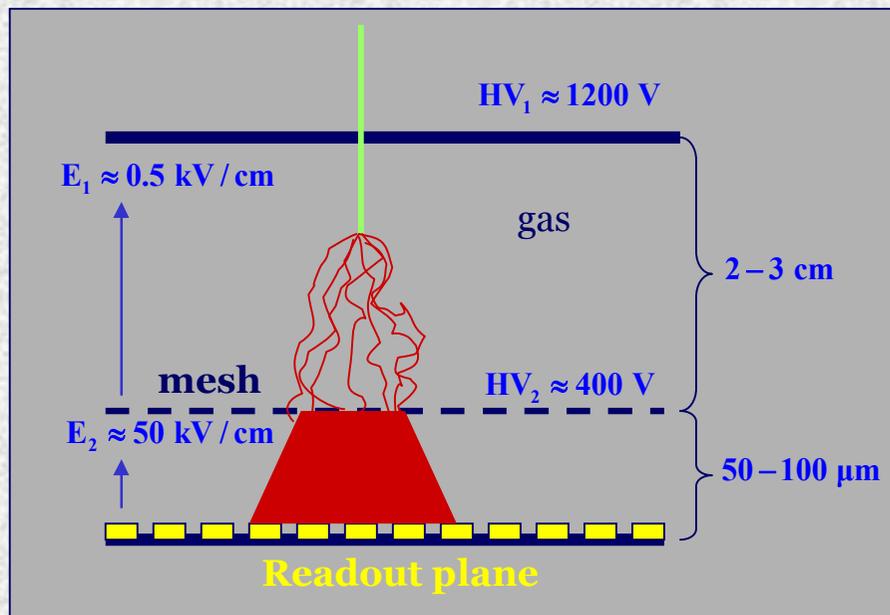
- Helioscopes.
- Laser experiments.
- Coherent Bragg scattering in crystals.
- Astrophysical arguments.
- Cosmological arguments.
- Dark matter experiments.



# MICROME GAS



# Operation principle

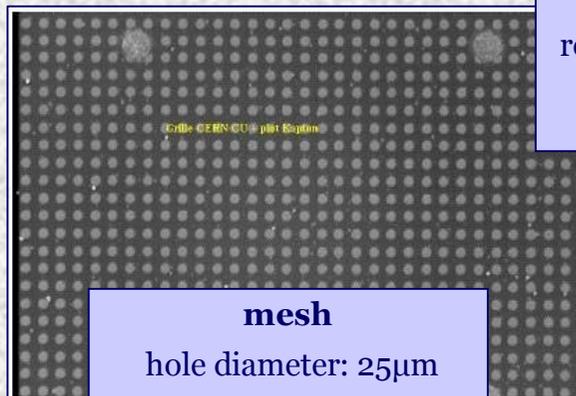
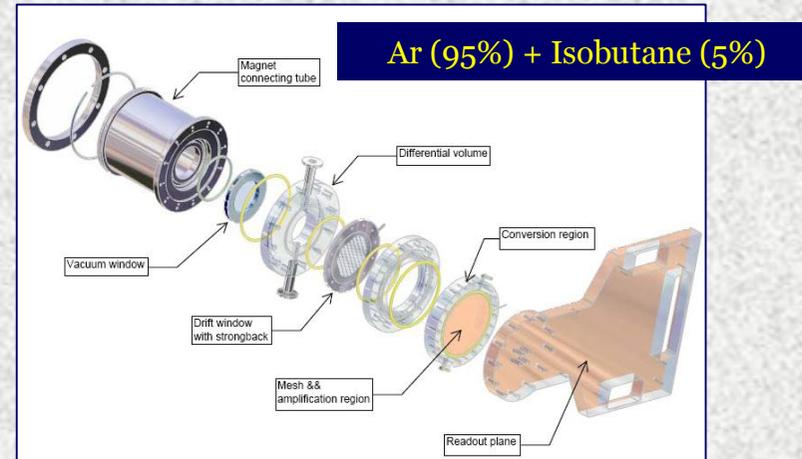




# The Micromegas of CAST

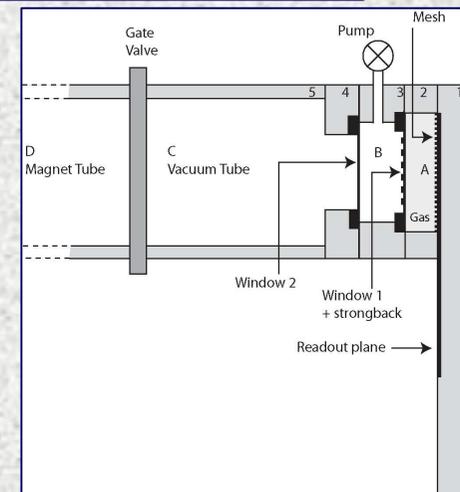
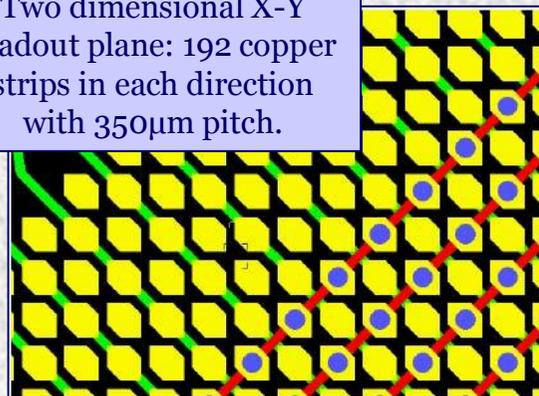
## Specifications

1. X-ray detection capability (energy 1-10 keV).
2. Typical energy resolution (20% at 5.9 keV).
3. Good efficiency.
4. Spatial sensitivity ~ 1 mm.
5. Calibration rate: 50Hz, data rate: 1Hz.



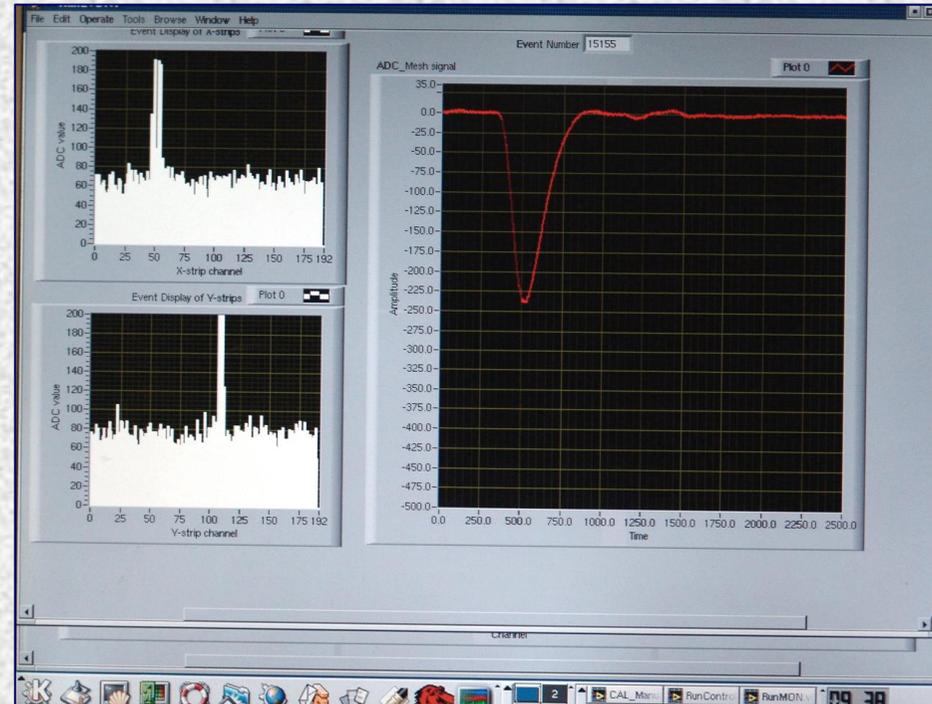
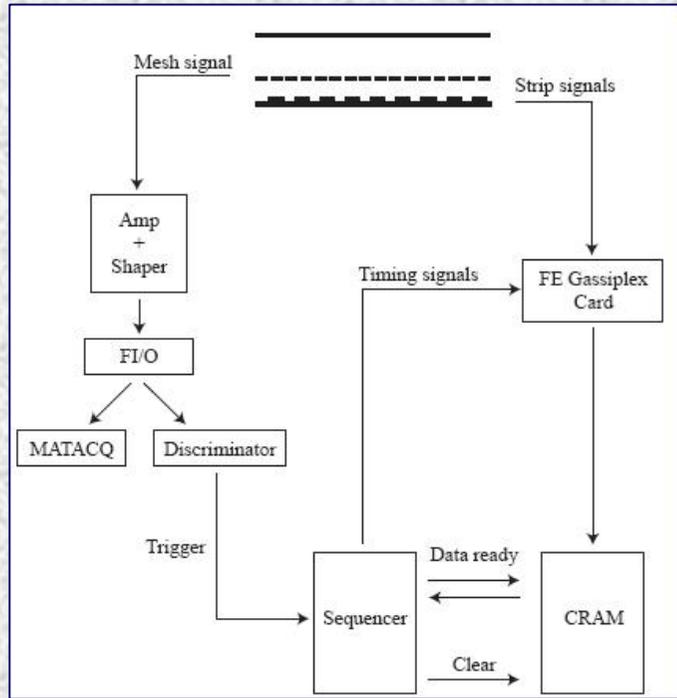
**mesh**  
hole diameter: 25μm  
pitch: 50μm

Two dimensional X-Y readout plane: 192 copper strips in each direction with 350μm pitch.





# Electronics



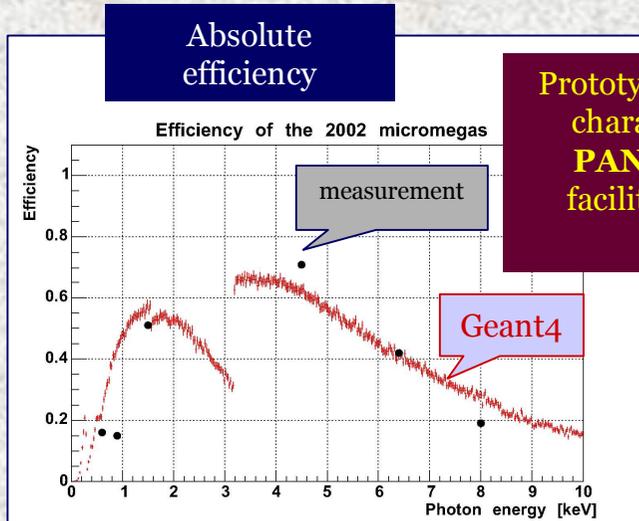
## Data Acquisition System (DAQ)

Written in LabView (Linux+Windows), developed in NCSR  
"Demokritos" (T. Geralis, G. Fanourakis).

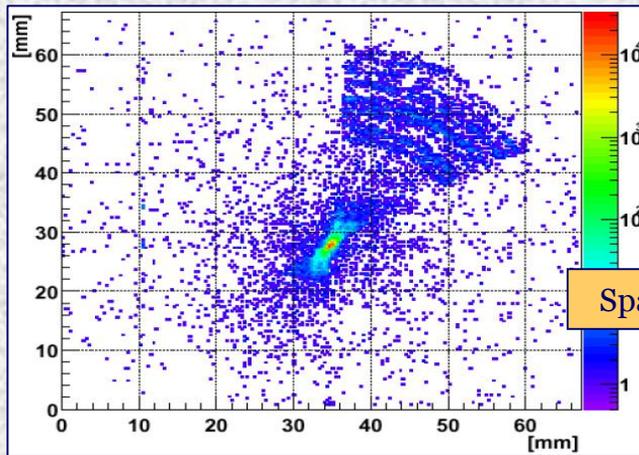
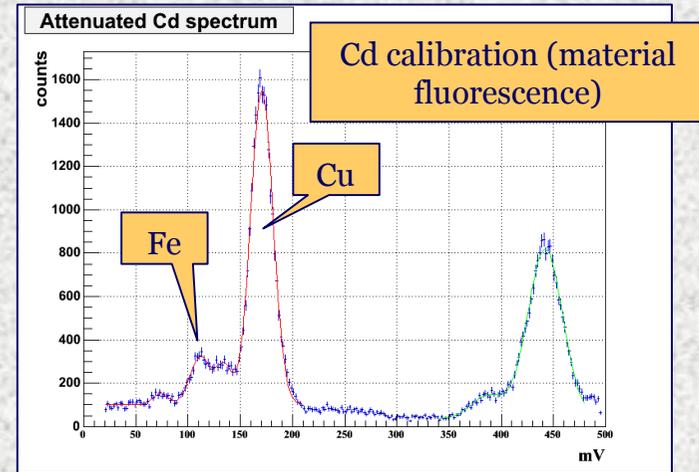
Operation: **Non Paralyzable**. Dead time: **14 msec**.



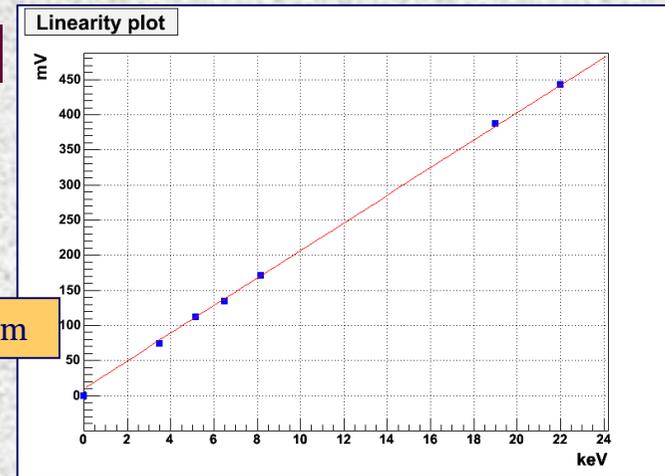
# Characterization



Prototype Micromegas:  
characterization at  
**PANTHER, MPE**  
facilities – Munich,  
2002.



**Linearity check**





# ANALYSIS & RESULTS



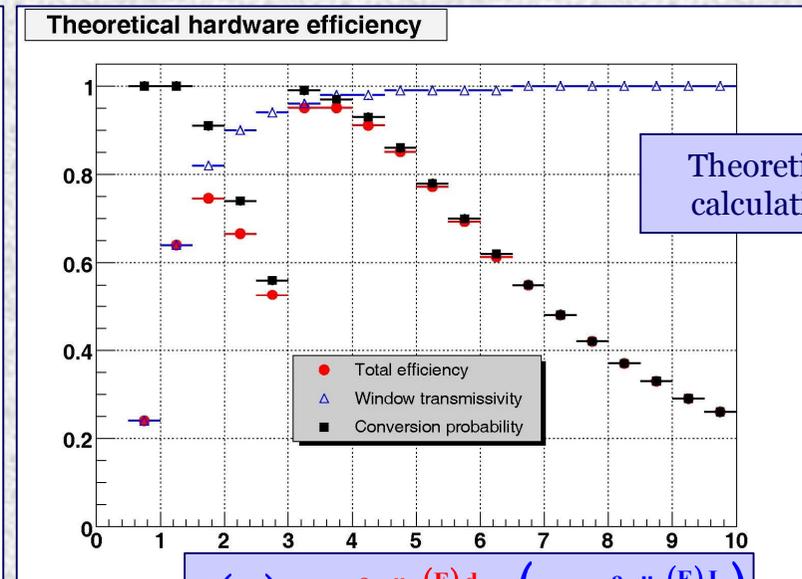
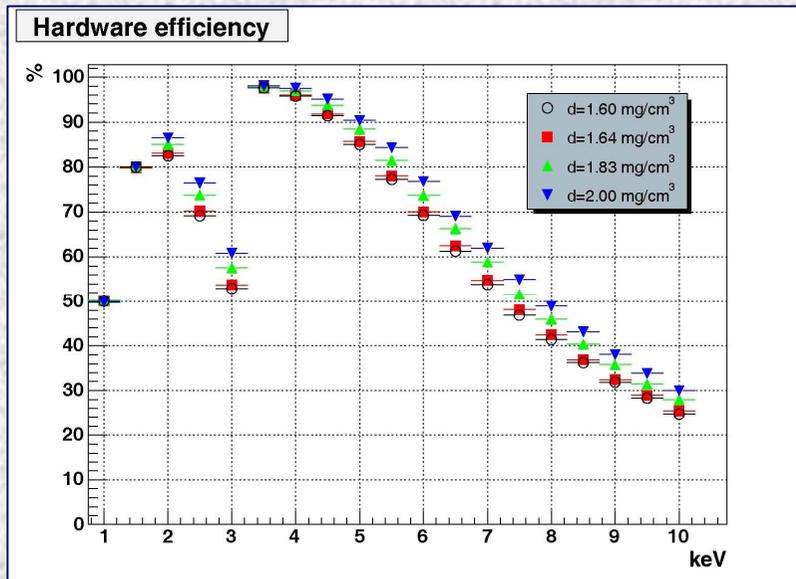
# Simulation (introduction)

## Description

- GEANT4 toolkit.
- Exact reconstruction of the detector's geometry and materials.
- Goals:
  1. Calculation of the detector efficiency for photons hitting transversely the side facing the magnet bore.
  2. Influence of gas density variation in the detector.
  3. Qualitative reproduction of the measured background.
- The simulation takes into account the energy deposition in the detector and NOT the full event reconstruction through the whole detection chain.



# Simulation (efficiency)



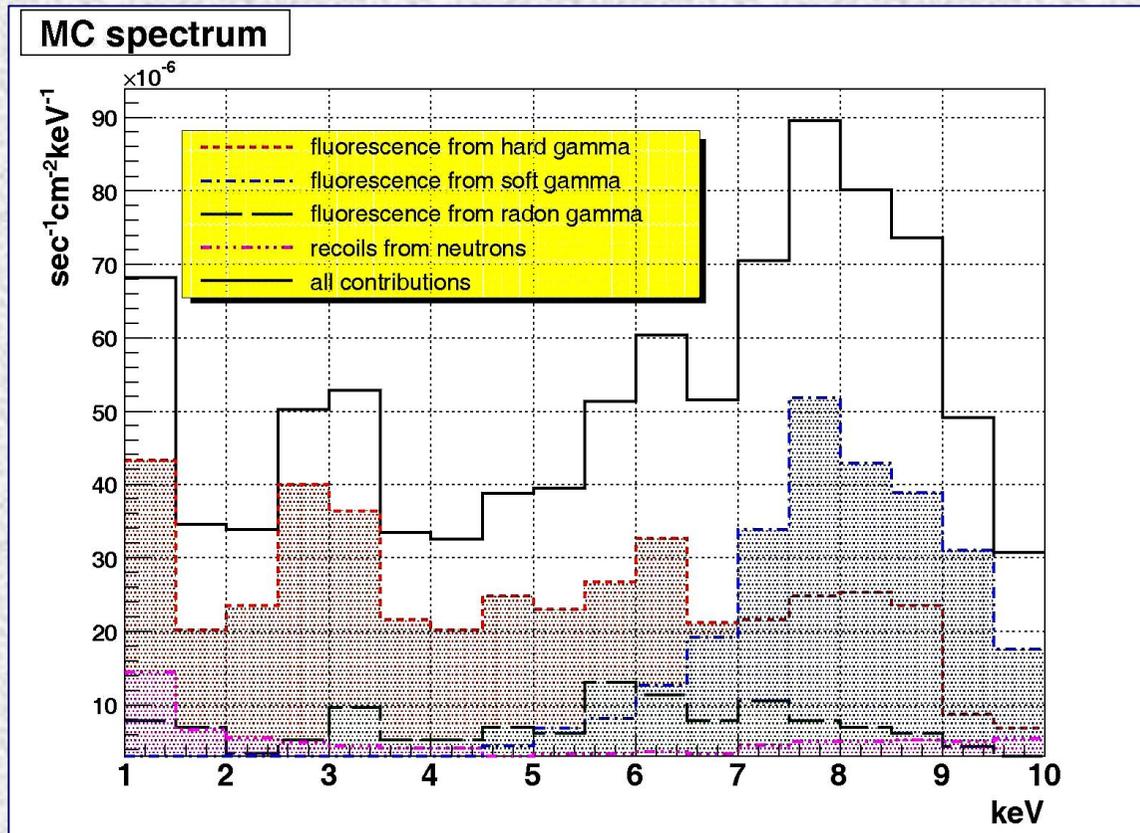
Full agreement between simulation and theoretical calculation.

$$\alpha(E) = e^{-\rho_w \cdot \mu_w(E) \cdot d} \times (1 - e^{-\rho_g \cdot \mu_g(E) \cdot L})$$

↑ windows      ↑ gas



# Simulation (background)



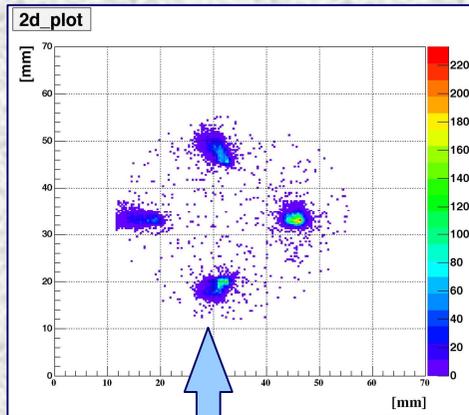
- Photons and neutrons of various energies are thrown to the detector isotropically.
- Event counting by local energy deposition.
- Normalization is arbitrary (based on the reconstruction of the Cu peak).
- Energy resolution is determined by the calibration measurements:

$$\delta = \sqrt{\frac{6(\text{keV})}{E}} \times 20\%$$

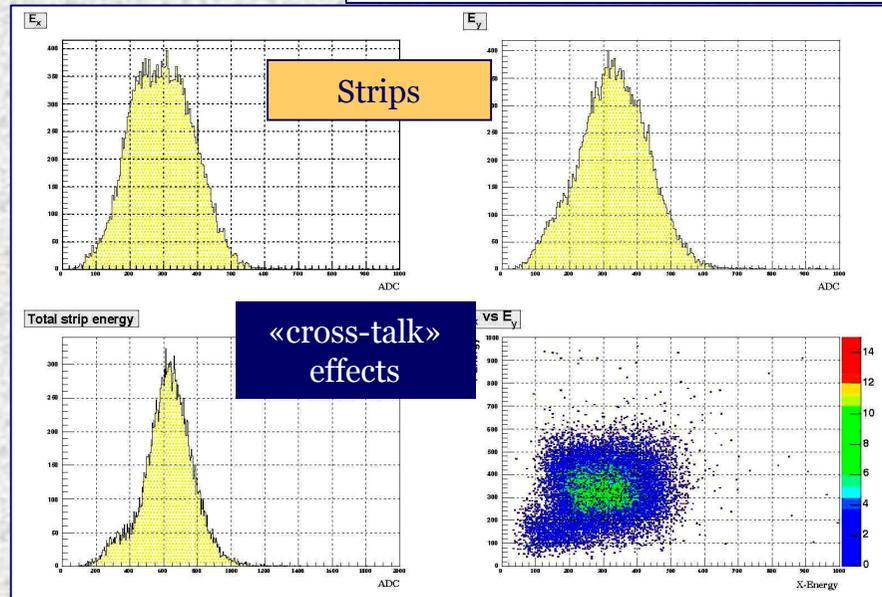
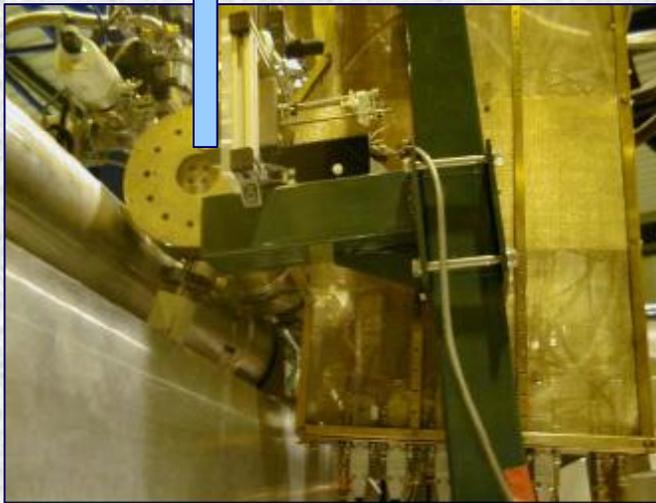
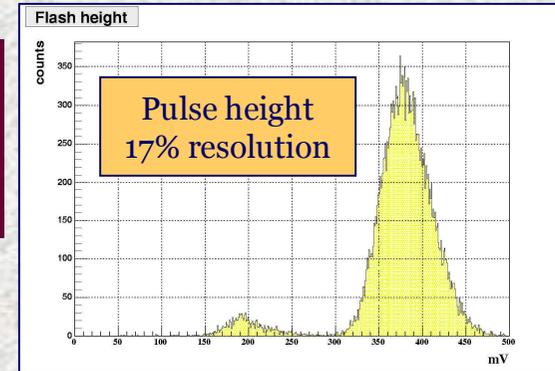
The background spectrum is determined by the fluorescence of the detector's building materials.



# 2003 - Calibration

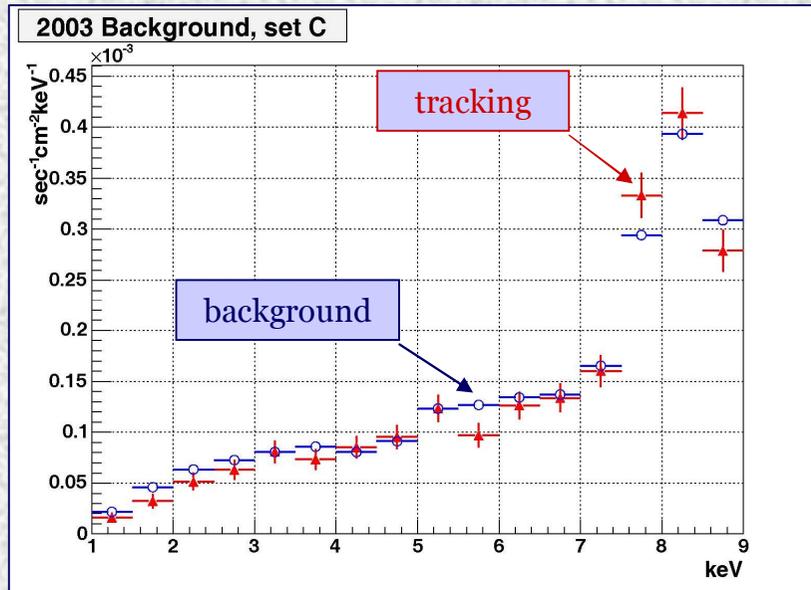


Prototype Micromegas V3  
Conversion gap: 18 mm  
Amplification gap: 50  $\mu\text{m}$





# 2003 - Data



- Event selection criteria**
1. Single cluster.
  2. Fiducial cut: circle with area  $A = 15.2 \text{ cm}^2$ . Magnet bore =  $14.5 \text{ cm}^2$ .
  3. Pulse risetime, width and ratio height/integral at most 2 standard deviations from the mean value.

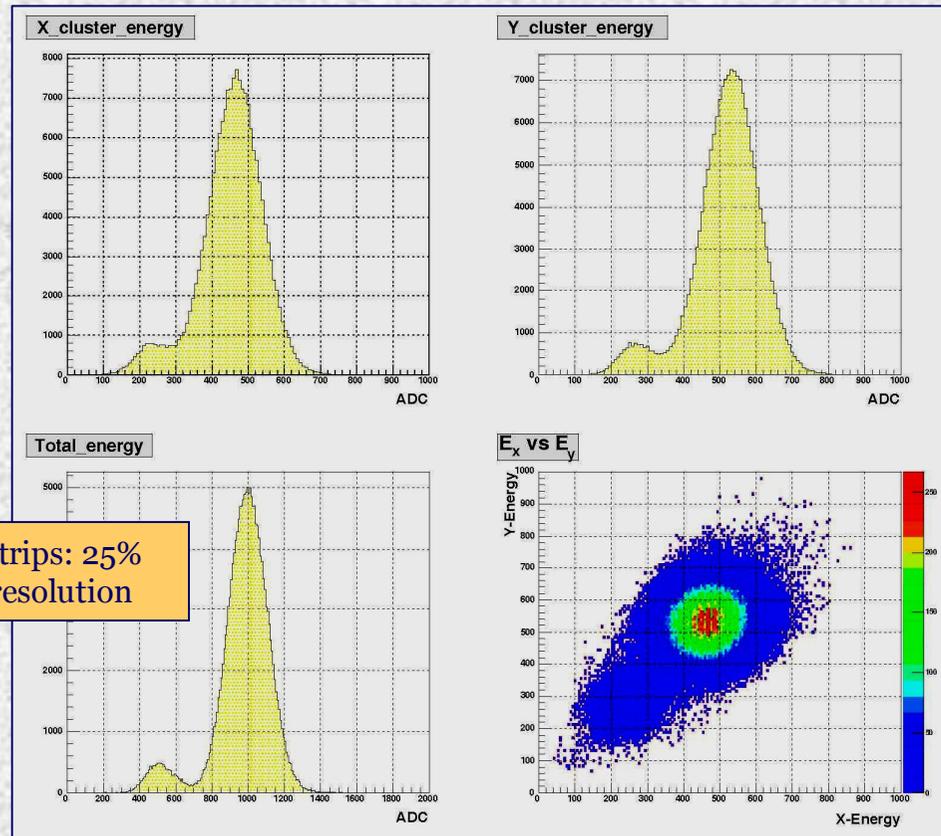
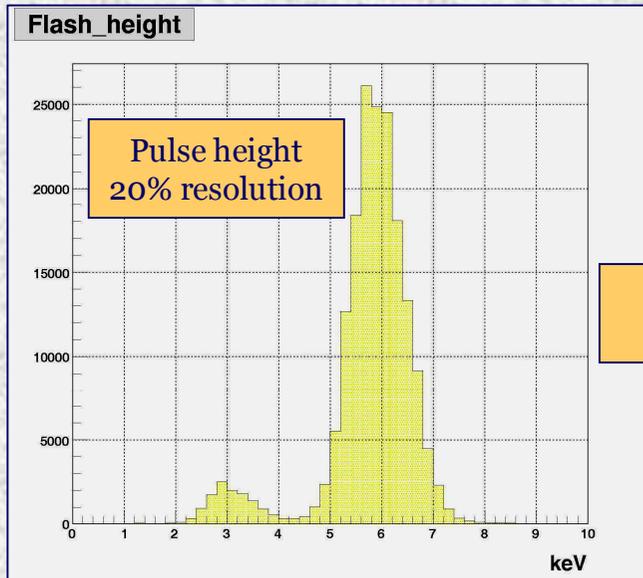
Three different datasets due to different timing of the pulses.

SET	Tracking time (h)	Background time (h)	$r_t$ $10^{-5}\text{s}^{-1}\text{cm}^{-2}\text{keV}^{-1}$	$r_b$ $10^{-5}\text{s}^{-1}\text{cm}^{-2}\text{keV}^{-1}$	Energy range (keV)
A	56.5	592	$8.4 \pm 0.2$	$8.73 \pm 0.06$	1 - 7.5
B	11.6	143	$10.7 \pm 0.5$	$11.35 \pm 0.14$	1 - 8.0
C	24.9	477	$13.5 \pm 0.3$	$13.91 \pm 0.08$	1 - 9.0



# 2004 - Calibration

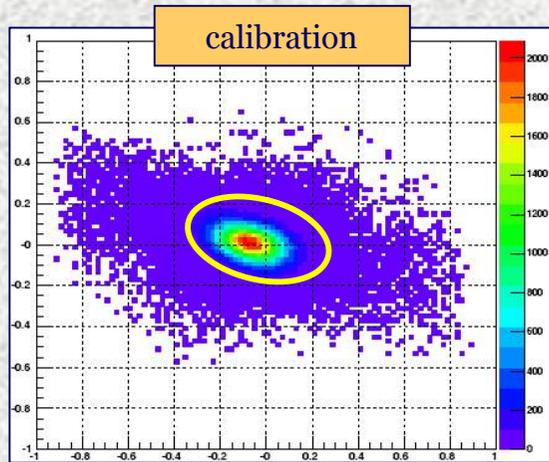
Prototype Micromegas V4  
Conversion gap: 23.5 mm  
Amplification gap: 100  $\mu\text{m}$



Strips: 25%  
resolution



# Multivariate analysis (general)



Goal:  
Maximum efficiency and purity.

$N$  variables  $x_i$  with mean values  $\mu_i$  following Gaussian distributions.

discriminant quantity

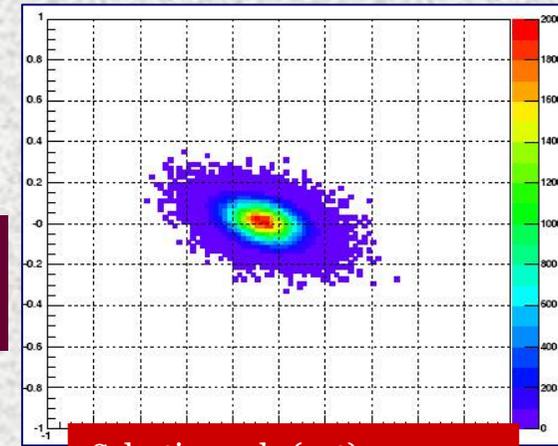
$$q = \sum_{ij} (x_i - \mu_i) V_{ij}^{-1} (x_j - \mu_j)$$

Covariance matrix

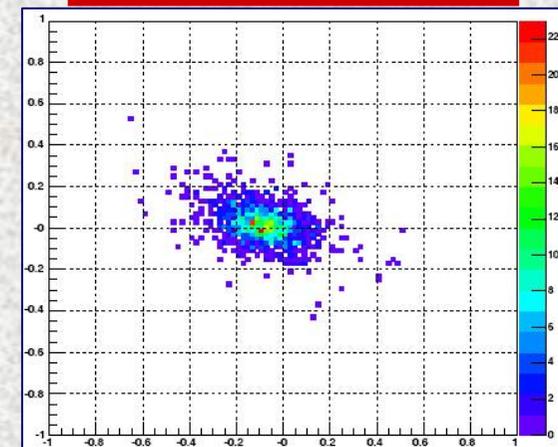
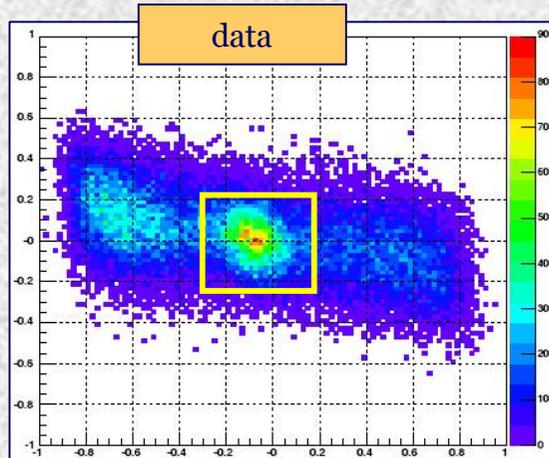
$$V_{ij} = \langle (x_i - \mu_i)(x_j - \mu_j) \rangle$$

$\chi^2_N$  distribution

$$f_N(q) = \frac{1}{2^{N/2} \Gamma(N/2)} \cdot q^{N/2-1} \cdot e^{-q/2}$$



Selection rule (cut):  $q < q_{\max}$



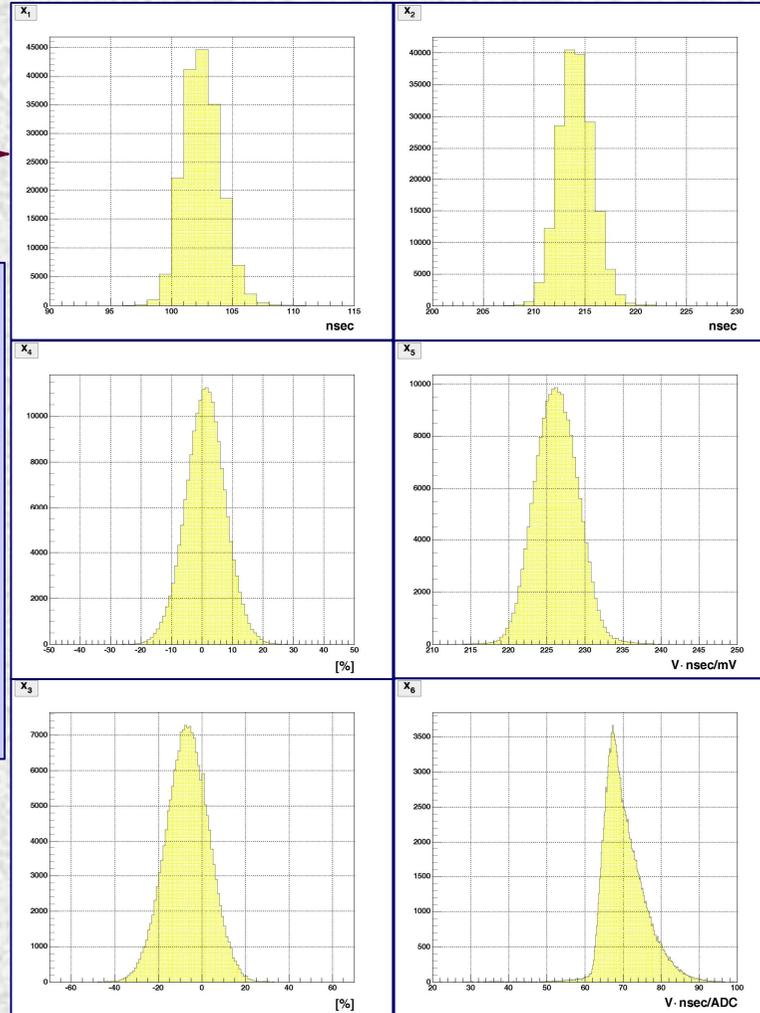


# Multivariate analysis (discriminant quantities)



**calibration** →

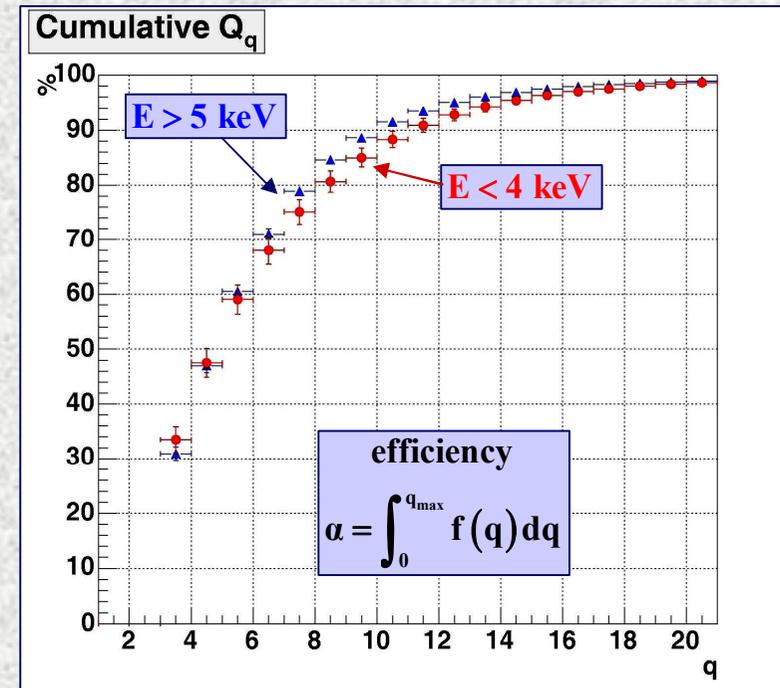
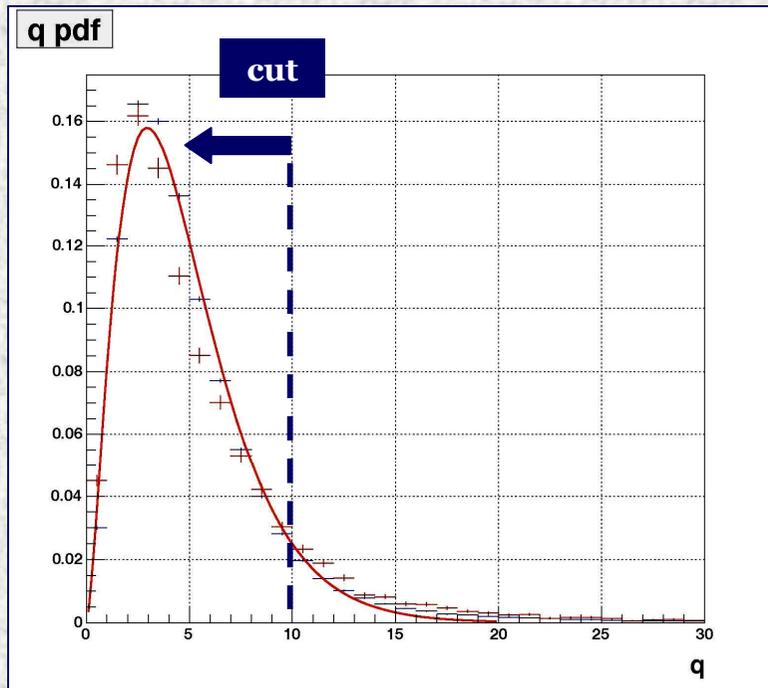
$x_1 = \text{risetime}$   
 $x_2 = \text{width}$   
 $x_3 = \frac{e_x - e_y}{e_x + e_y}$  (charge asymmetry)  
 $x_4 = \frac{\sigma_y m_x - \sigma_x m_y}{\sigma_y m_x + \sigma_x m_y}$  (multiplicity asymmetry)  
 $x_5 = \frac{I}{10^3 H}$  (height-integral correlation)  
 $x_6 = \frac{I}{10^3 (e_x + e_y)}$  (pulse-strip correlation)





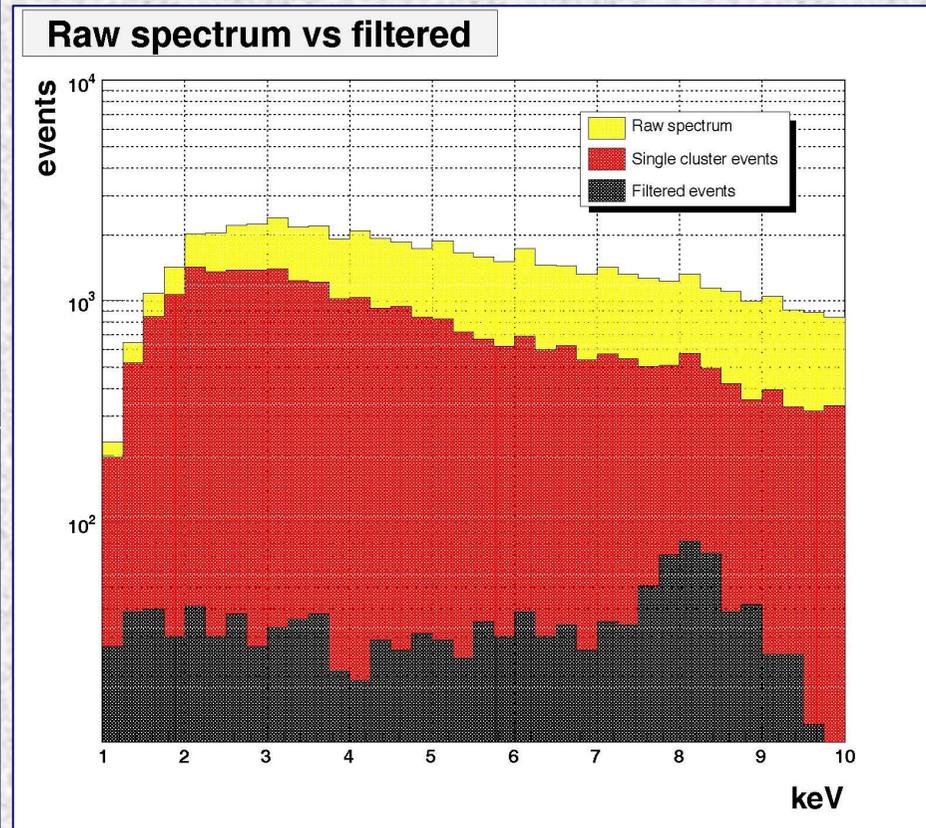
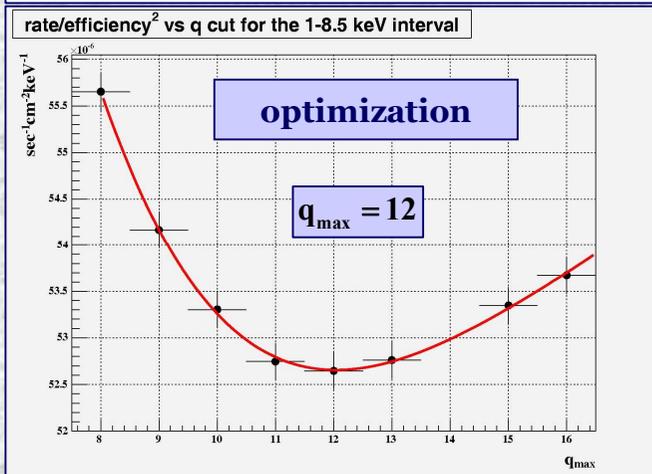
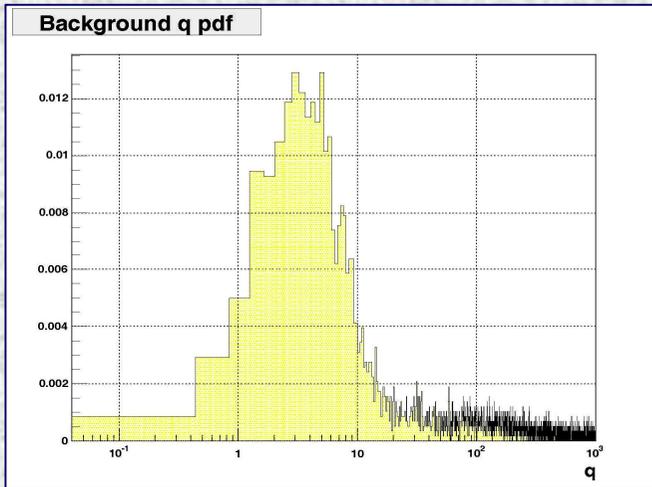
# Multivariate analysis (calibration)

$$\left. \begin{array}{l} \mu_i(E) = \text{const.} \\ \sigma_i(E) \sim E^{-3/4} \end{array} \right\} \Rightarrow V_{ij}^{-1}(E) = \left(\frac{E}{E_0}\right)^{3/2} V_{ij}^{-1}(E_0) \longrightarrow \tilde{q} = \left(\frac{E_0}{E}\right)^{3/2} \times \sum_{ij} (x_i - \mu_i) \cdot V_{ij}^{-1}(E_0) \cdot (x_j - \mu_j)$$



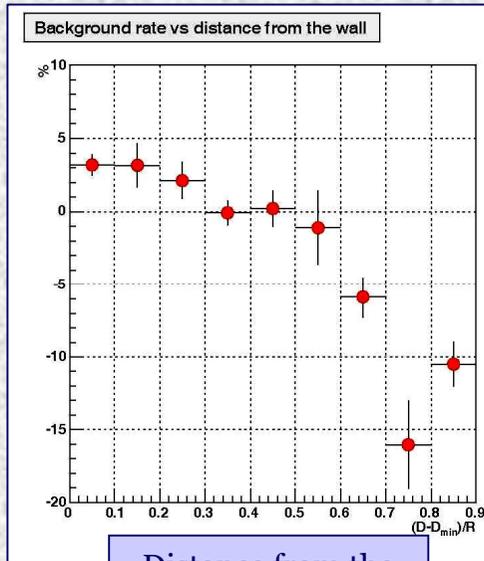


# Multivariate analysis (data)

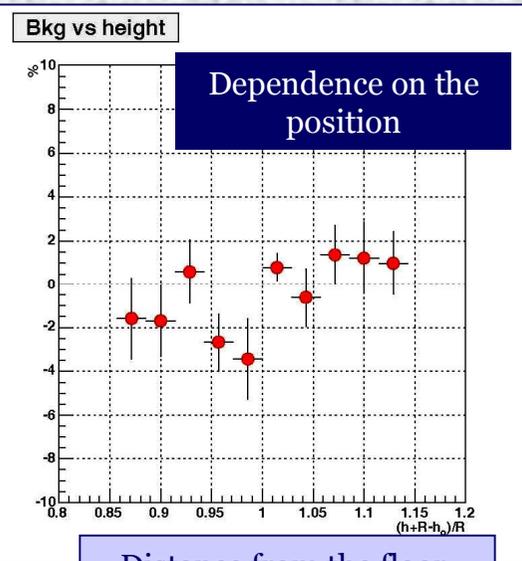




# Background study



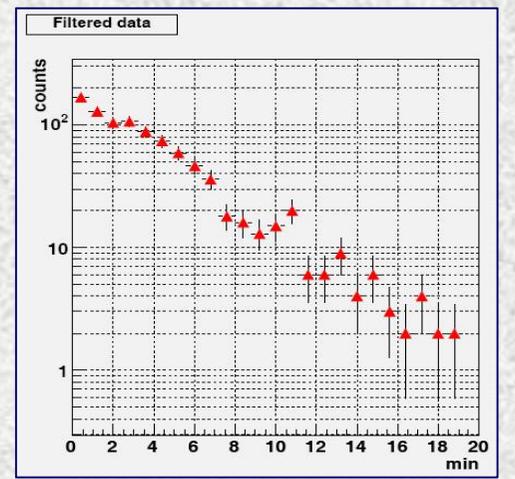
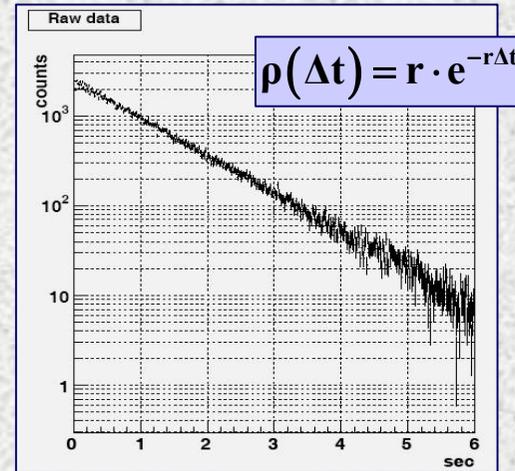
Distance from the closest wall.



Distance from the floor.

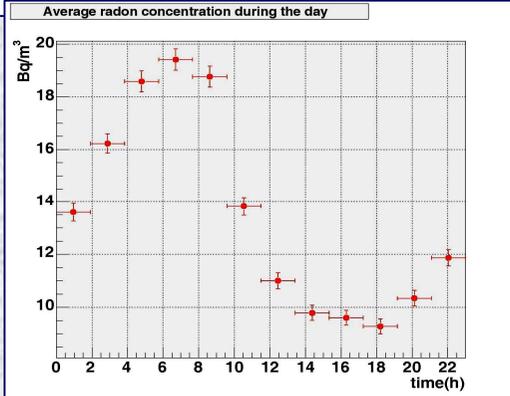
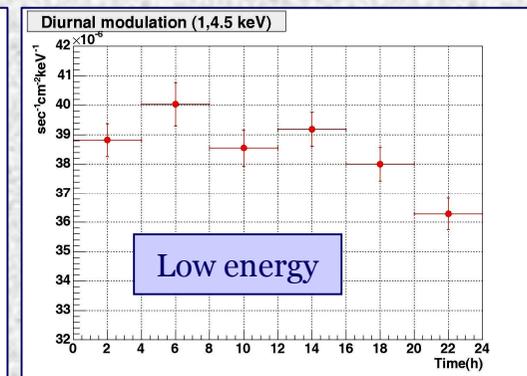
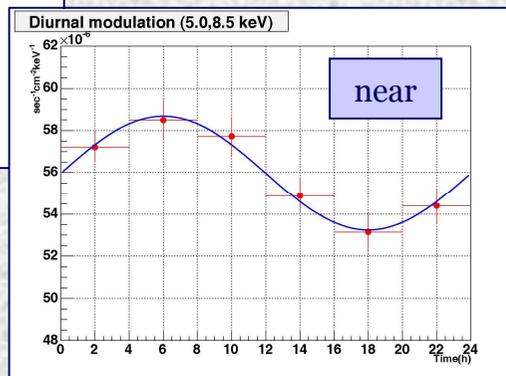
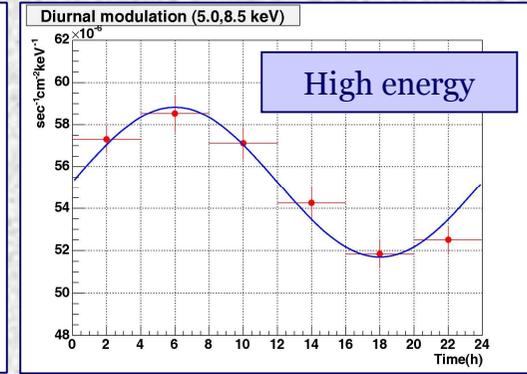
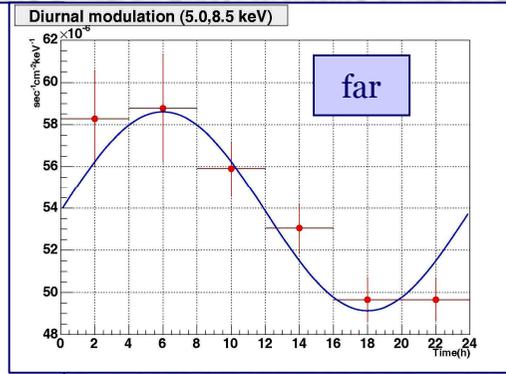
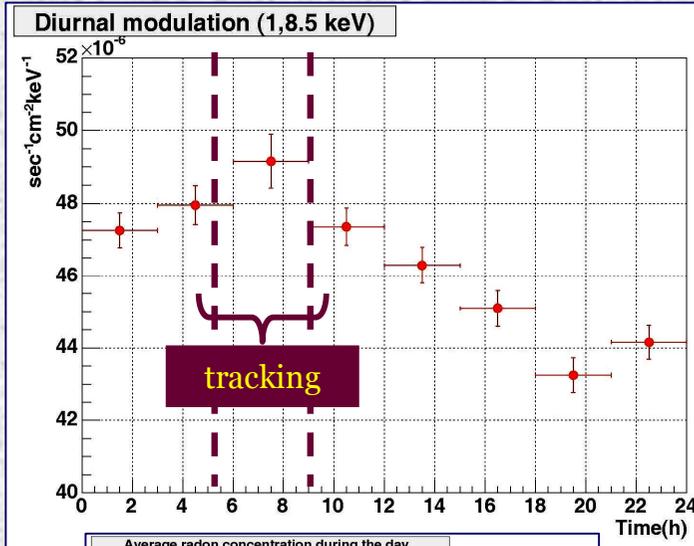
Dependence on the position

- $\Delta E = 1 - 8.5$  keV.
- $A = 15.2$  cm<sup>2</sup>.
- integration time: > 3000 h.
- trigger rate: ~ 1 Hz.
- filtered rate: ~ 18 events/h.
- cut efficiency: 92%.





# Daily variation



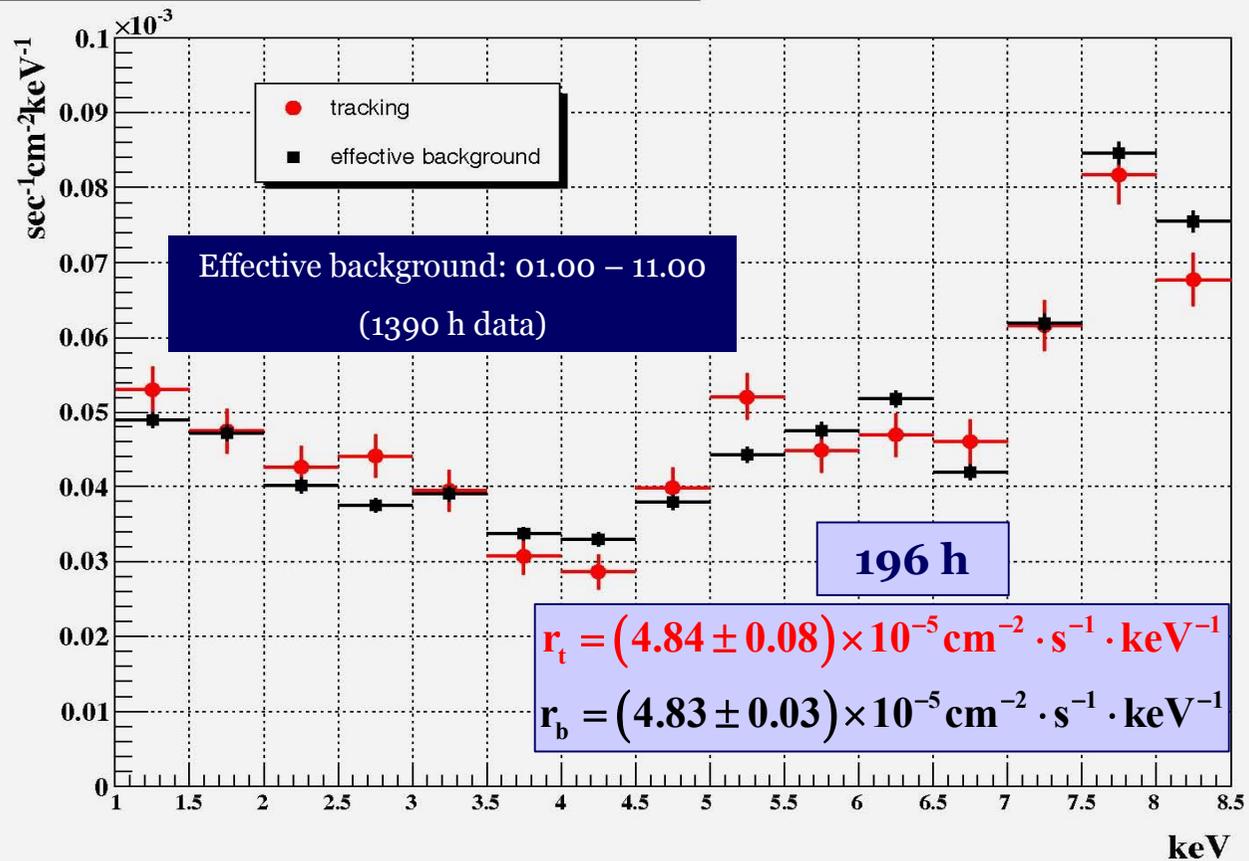
Clear daily variation, probably due to the variation of the Rn concentration.



# Solar data

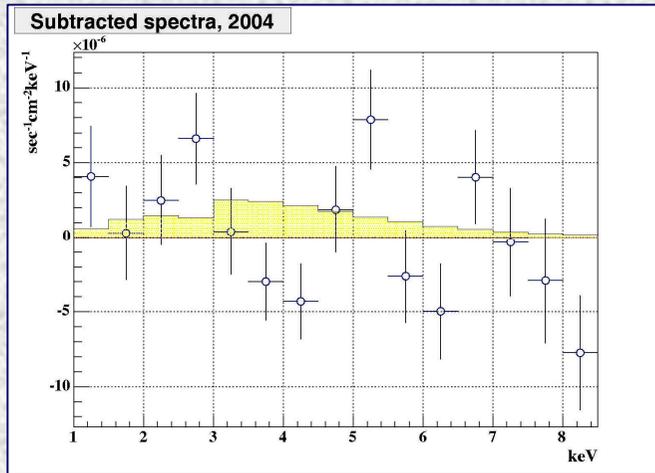


Tracking vs Background, 2004





# Calculation of an upper limit for $g_{\alpha\gamma\gamma}$



$$N_i^a = f(E_i) \times A_m \times t \times P_{a \rightarrow \gamma} \times \alpha_h(E_i) \times \alpha_s \times dE$$

$$N_i^m = (r_i^t - r_i^b) \times t \times dE \times A$$

$$\chi^2 = \sum_{\text{bins}} \left( \frac{N_i^a - N_i^m}{\sigma_i} \right)^2$$

$$\sigma_i = \sqrt{N_i^t + \left( \frac{t_t}{t_b} \right)^2 N_i^b}$$

fitting parameter:  $C = g_{10}^4$

$f(E) \sim g_{10}^2$  (axion flux)

$g_{10} = g_{\alpha\gamma\gamma} \times (10^{10} \text{ GeV})$

$A_m = 14.5 \text{ cm}^2$

$t = 196 \text{ h}$

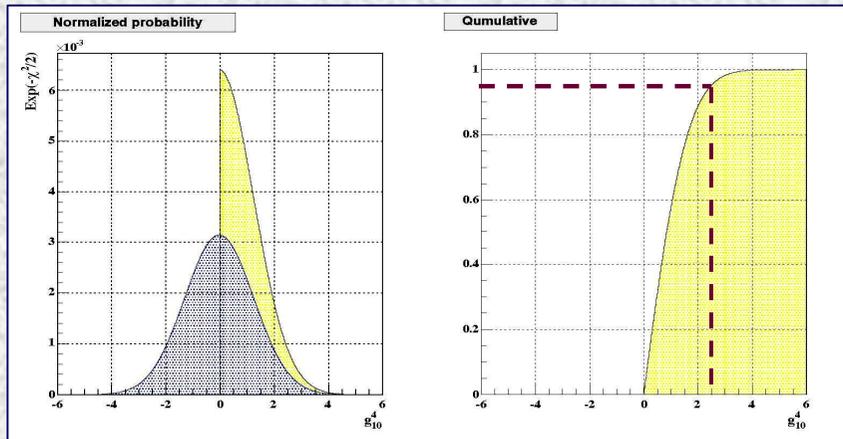
$P_{a \rightarrow \gamma} = 0.25 \times B^2 L^2 g_{\alpha\gamma\gamma}^2$ ,  $m_a < 0.02 \text{ eV}$

$\alpha_h(E)$ : hardware efficiency

(detection+strongback+dead time)

$\alpha_s$ : software efficiency

$dE = 0.5 \text{ keV}$



Bayes method

$$(1-a) \int_0^{+\infty} \exp\left(-\frac{\chi^2}{2}\right) dC \quad \int_0^{C_{\text{up}}} \exp\left(-\frac{\chi^2}{2}\right) dC$$

Degree of belief

Upper limit



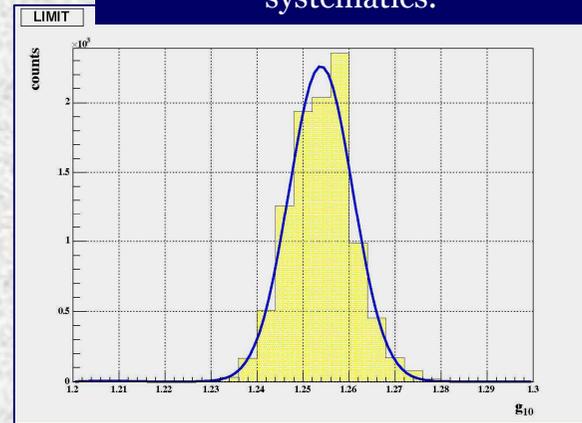
# Calculation of an upper limit for $g_{\alpha\gamma}$



1 - 8.5 keV

Effective background	Best fit	$\chi^2_{\min} / \text{ndf}$	$\chi^2_{\text{null}} / \text{ndf}$	95% CL $10^{-10} \text{ GeV}^{-1}$
00.00-12.00	$0.21 \pm 0.89$	24.73 / 14	24.75 / 15	1.27
01.00-11.00	$-0.03 \pm 0.9$	26.84 / 14	26.84 / 15	1.25
02.00-10.00	$-0.35 \pm 0.92$	30.11 / 14	30.18 / 15	1.23
03.00-09.00	$-0.67 \pm 0.94$	31.77 / 14	32.02 / 15	1.22

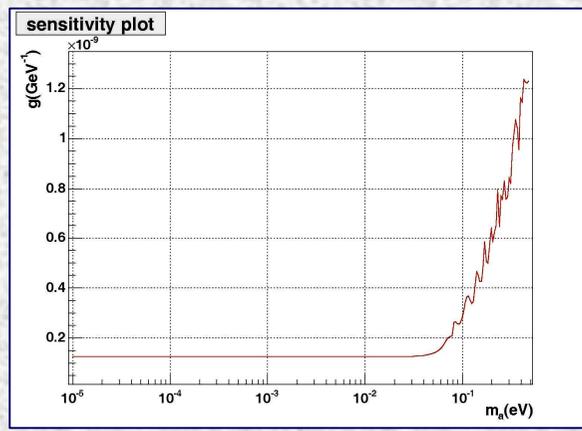
Simulation of the efficiency systematics.



1 - 7.5 keV

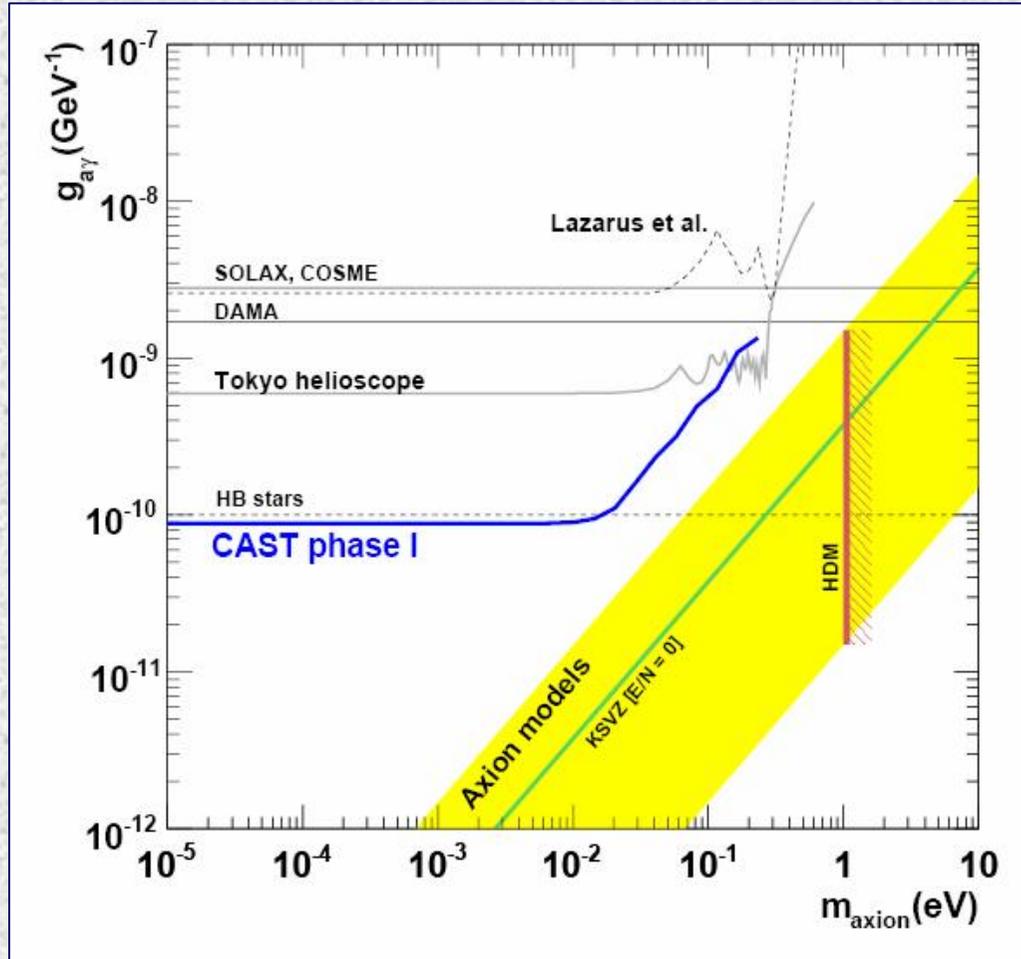
$g_{\alpha\gamma} < (1.25 \pm 0.02 \pm 0.007) \times 10^{-10} \text{ GeV}^{-1}$   
 $m_\alpha < 0.02 \text{ eV}$

Effective background	Best fit	$\chi^2_{\min} / \text{ndf}$	$\chi^2_{\text{null}} / \text{ndf}$	95% CL $10^{-10} \text{ GeV}^{-1}$
00.00-12.00	$0.29 \pm 0.9$	20.83 / 12	20.88 / 13	1.28
01.00-11.00	$0.06 \pm 0.9$	22.26 / 12	22.26 / 13	1.26
02.00-10.00	$-0.24 \pm 0.92$	23.93 / 12	23.97 / 13	1.24
03.00-09.00	$-0.55 \pm 0.94$	24.54 / 12	24.71 / 13	1.23





# CAST results



95% C.L. (Bayes) upper limit

$$g_{\alpha\gamma\gamma} < 8.8 \times 10^{-11} \text{ GeV}^{-1}$$

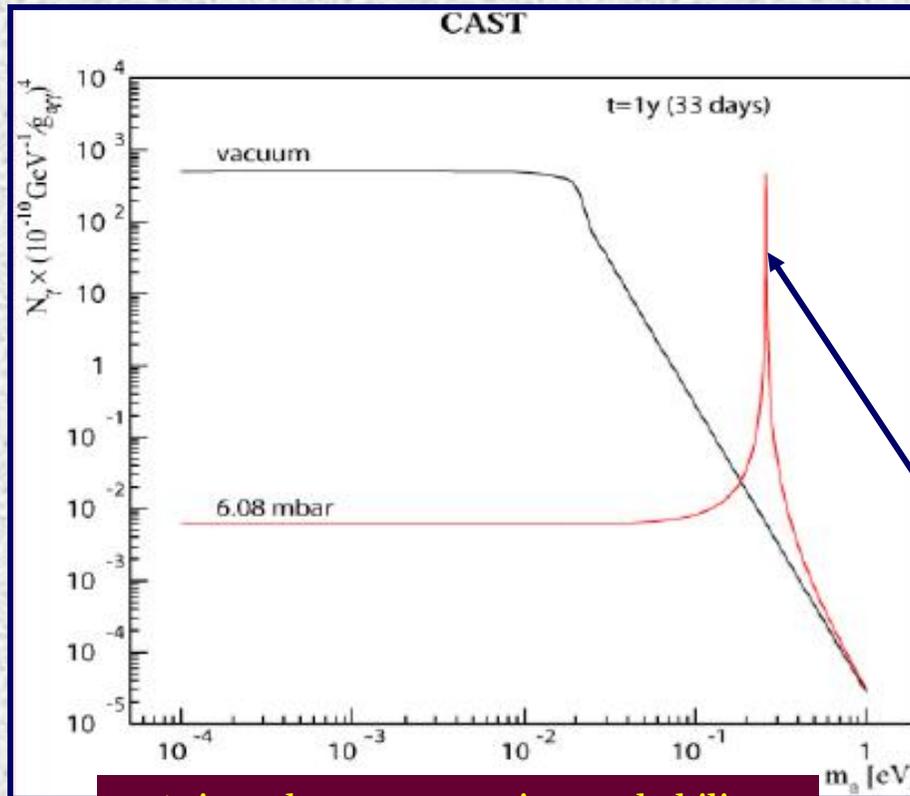
Submitted for  
publication (JCAP)

hep-ex/0702006 :

“An improved limit on the  
axion – photon coupling from  
the CAST experiment”



# CAST Phase II principle



Axion-photon conversion probability

$m_\gamma$  = effective photon mass

$$m_\gamma \approx \sqrt{\frac{4\pi\alpha N_e}{m_e}} = 28.9 \sqrt{\frac{Z}{A} \rho} \text{ eV}$$

$\rho$  = gas density (g/cm<sup>3</sup>)

$$q = \frac{|m_\alpha^2 - m_\gamma^2|}{2E} \text{ (momentum transfer)}$$

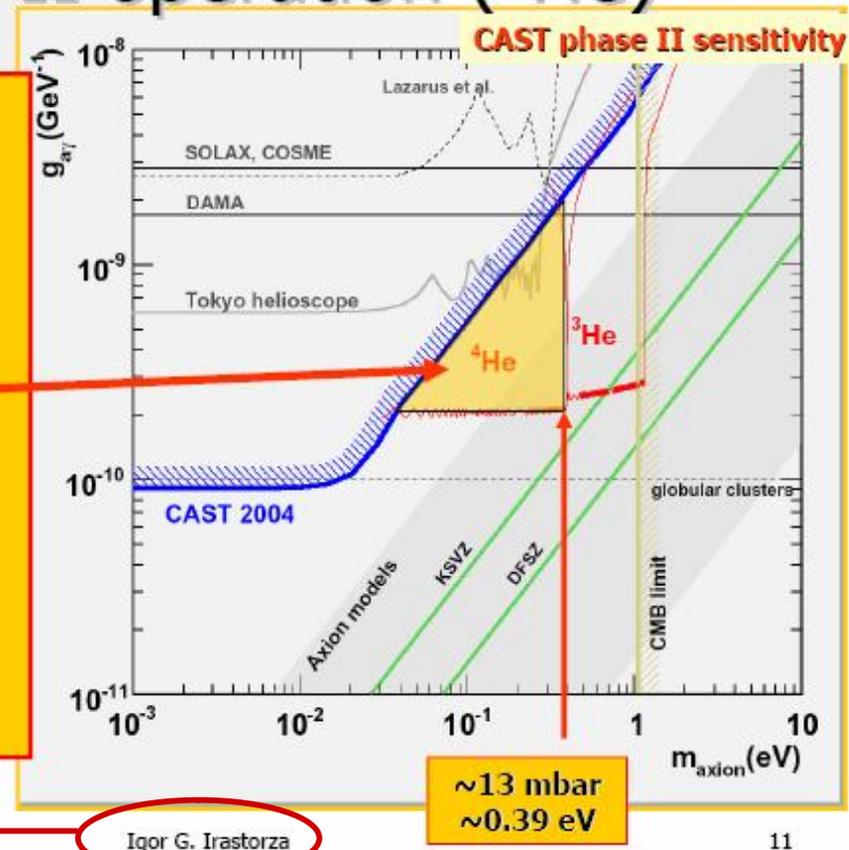
Coherence condition:  
 $qL \ll 1$



# CAST Phase II status & prospects

## Phase II operation ( $^4\text{He}$ )

- Data taking with  $^4\text{He}$  performed all along 2006
- $\sim 160$  density steps performed, reaching  $\sim 13$  mbar ( $\sim 0.39$  eV)
- Approximate explored region shown in plot
  - QCD axion models region is entered !!.
- Data now under analysis
- $^3\text{He}$  phase will start mid 2007 (upgrade works ongoing)



Moriond EW 2007

Igor G. Irastorza

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# Conclusions

- The Strong CP problem remains unresolved for 30 years and the dynamic Peccei-Quinn solution is the most viable. Axion searches are too important to ignore.
  - Following the analysis of the data collected with the Micromegas detector during CAST Phase I **there is no axion signal.**
  - The missing axion signal leads to the calculation of an upper limit for the axion-photon coupling constant which improves the limits from previous experiments by a factor  $\sim 7$ .
  - CAST continues with Phase II, entering unexplored regions of the axion parametric space.
  - The Micromegas detector is a stable and reliable structure, as proved by its performance in the CAST experiment.
- 
- The 2003 data were used mostly for the understanding of the detector, while the 2004 data were used to make Physics analysis.
  - The spectrum of the measured background from the Micromegas detector is fully understood, with the aid of the simulation and it comes from the fluorescence of its constructing materials.
  - The multivariate analysis is more suitable when the event selection rules are based on correlated discriminant quantities.
  - The dominant systematical error is the diurnal variation of the measured background and it was taken into account by defining an effective background. The total effect of the systematical errors to the calculated upper limit of the axion-photon coupling constant is less than 2%.



# Presentations



## Publications

1. Performance of the Micromegas detector in the CAST experiment,  
**Nucl.Instrum.Meth.A573:38**, 2007
2. First results from the CERN Axion Solar Telescope (CAST),  
**Phys.Rev.Lett.94:121301**, 2005.
3. A low background Micromegas detector for axion searches,  
**Nucl.Instrum.Meth.A535:309**, 2004.
4. The use of the Micromegas technology for a new imaging system,  
**Nucl.Instrum.Meth.A527:62**, 2004.
5. An improved limit on the axion-photon coupling from the CAST experiment,  
**Submitted to JCAP.**
6. The Micromegas detector of the CAST experiment,  
**Submitted to New J. Phys.**

## Conferences

1. HEP2004, Chios.
2. PSD7 2005, Liverpool.
3. HEP2006, Ioannina.

## CAST collaboration meetings

1. 17<sup>th</sup> Col.Meeting (2003-CERN).
2. 23<sup>rd</sup> Col.Meeting (2004-Munich).
3. 24<sup>th</sup> Col.Meeting (2005-CERN).
4. 25<sup>th</sup> Col.Meeting (2005-CERN).
5. 27<sup>th</sup> Col.Meeting (2005-CERN).
6. 28<sup>th</sup> Col.Meeting (2006-Zaragoza).



# Acknowledgements

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I. G. Irastorza

T. Dafni

B. Lakić

NCSR  
"Demokritos"

Saclay

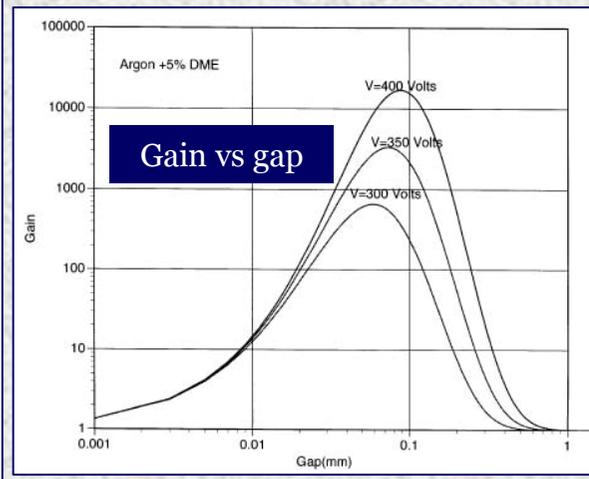
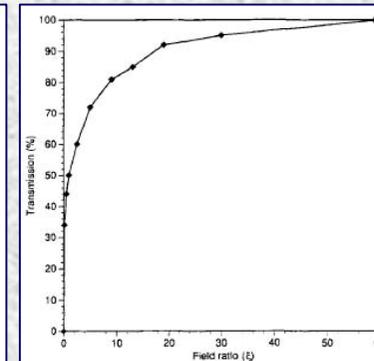
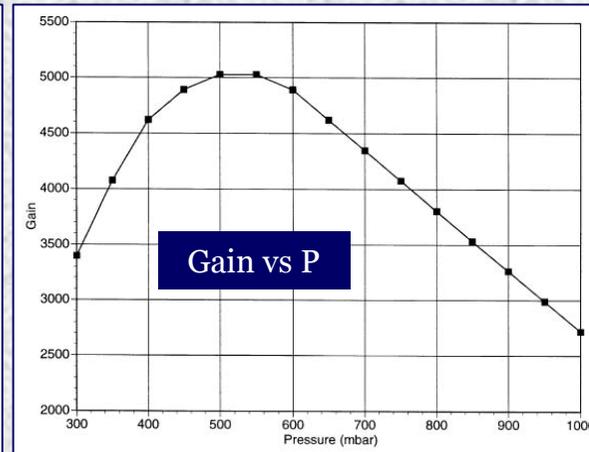
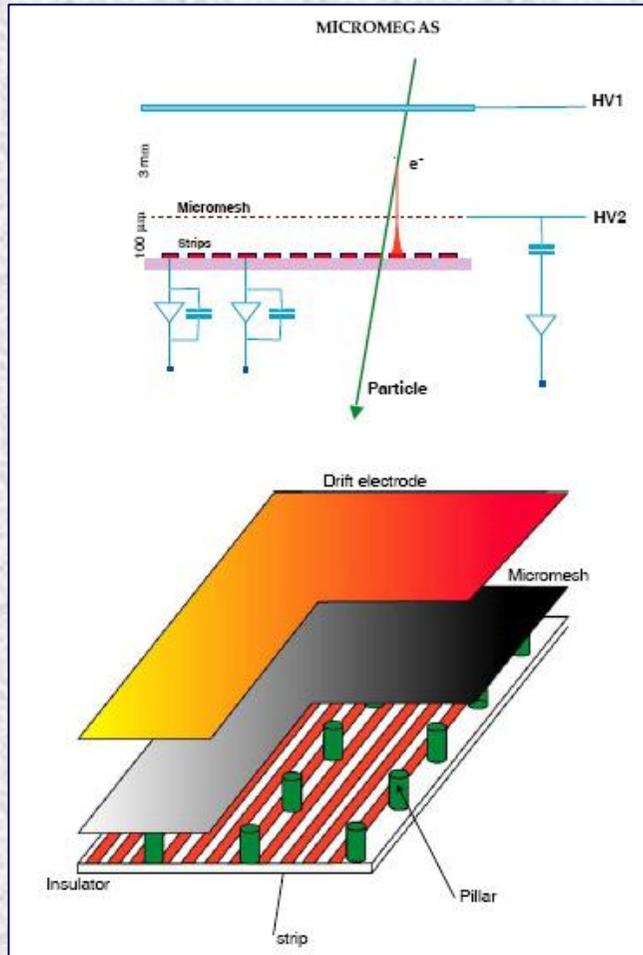
CAST



# BACKUP



# Micromegas



$$M = e^{\alpha \cdot d}$$

$$\alpha = p \cdot A \cdot \exp\left(-C \frac{pd}{V}\right)$$

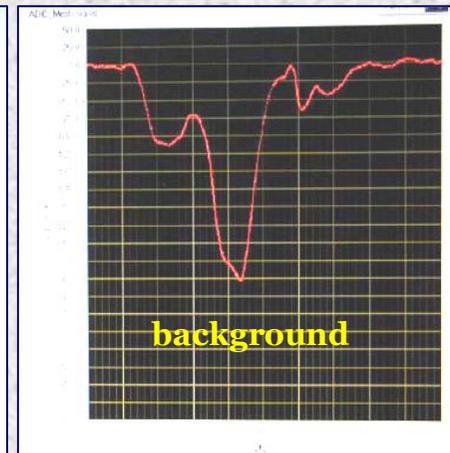
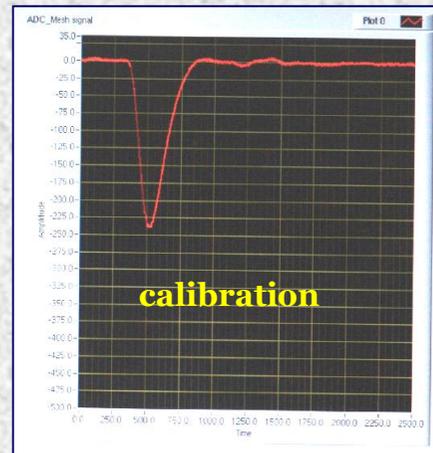
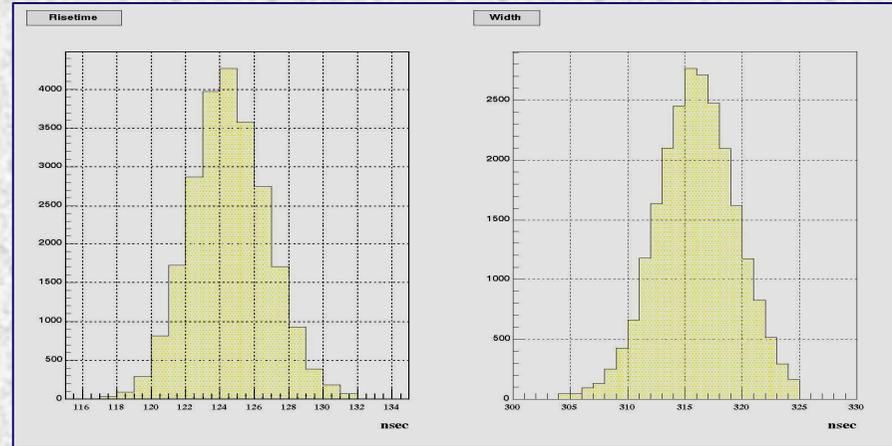
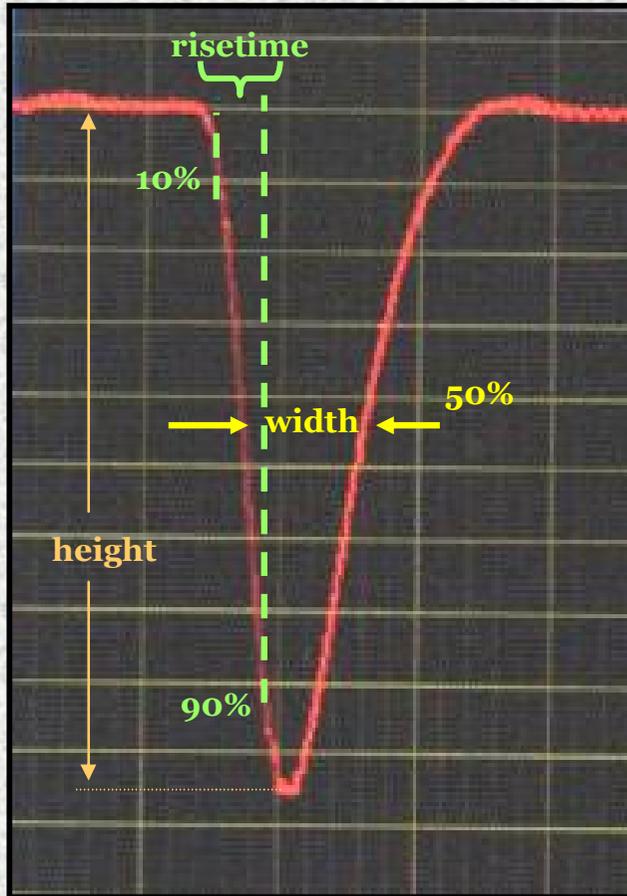
(Rose & Korff)

$$\frac{\delta M}{M} = \alpha \cdot d \cdot \left(1 - \frac{Cpd}{V}\right) \cdot \frac{\delta d}{d}$$

- $\alpha$ : Townsend coefficient.
- $d$ : amplification gap.
- $V$ : High voltage.
- $p$ : gas pressure.
- $A, C$ : gas constants.

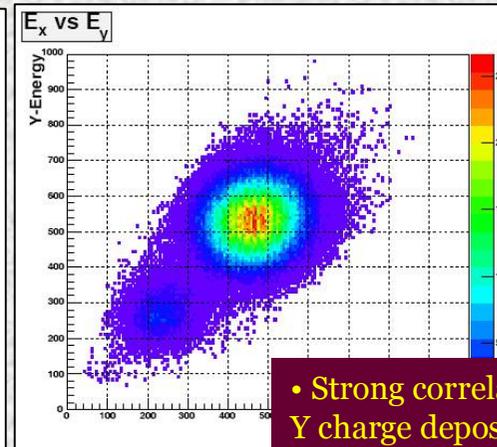
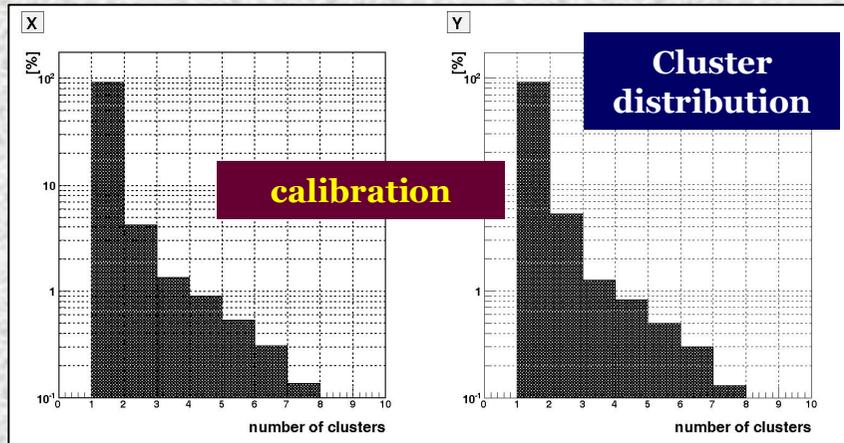


# Pulse analysis

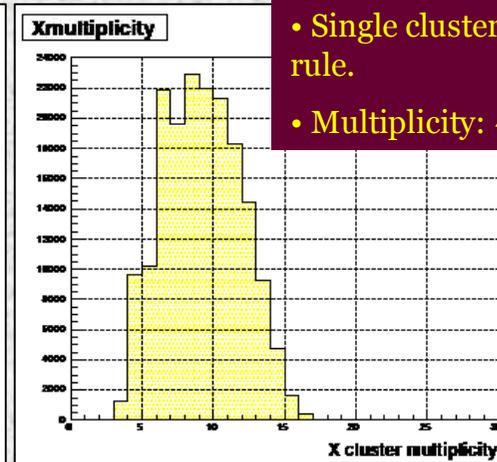
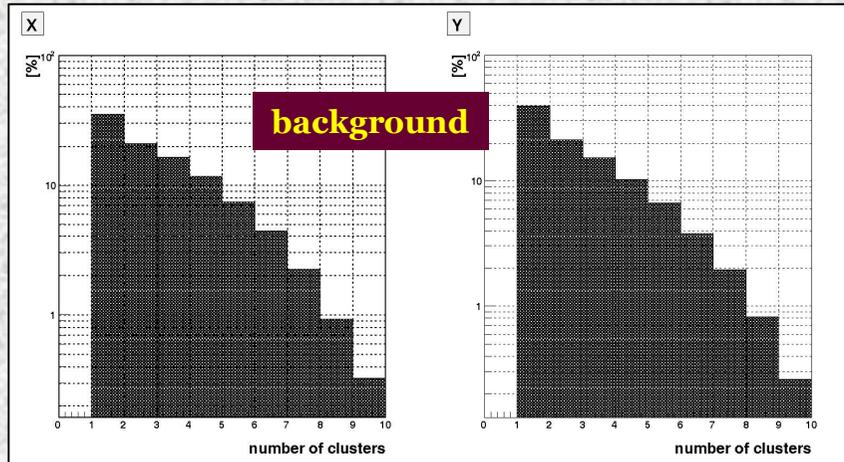




# Typical strip data



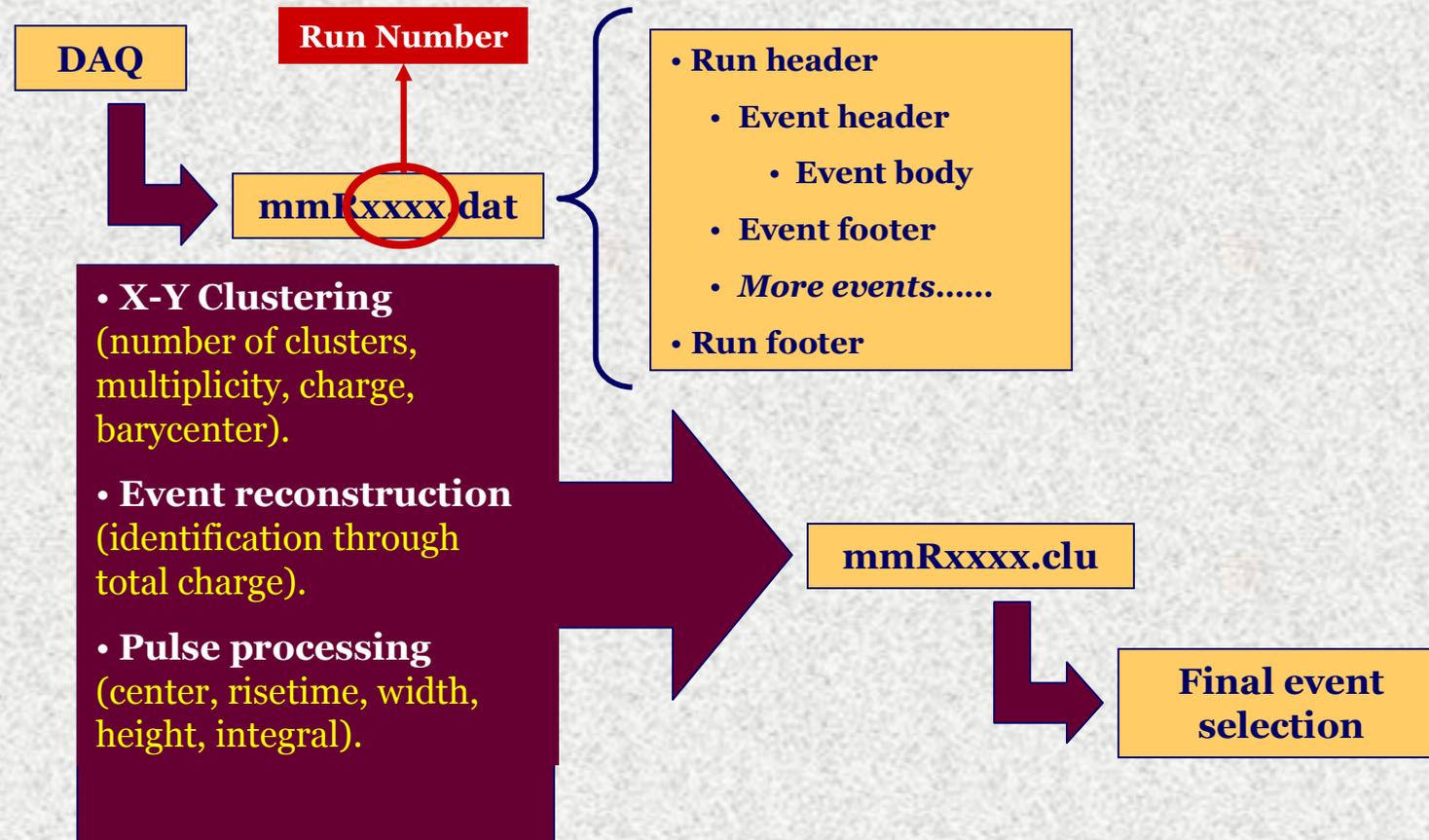
- Strong correlation between X and Y charge deposition.



- Single cluster: strong selection rule.
- Multiplicity: ~ 9 strips/cluster.



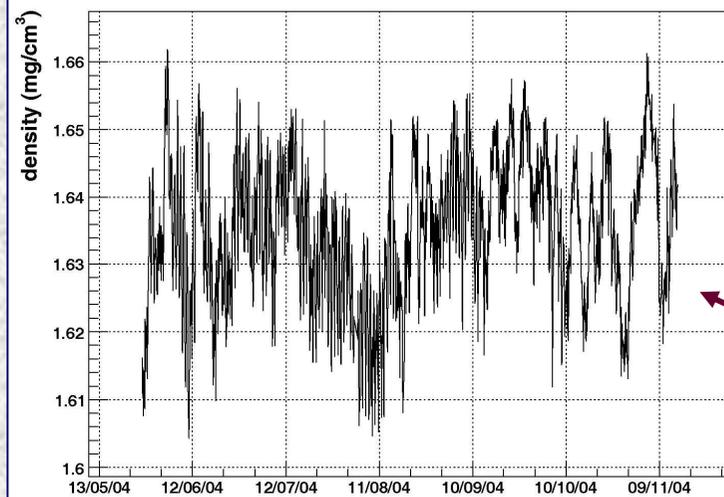
# Data processing



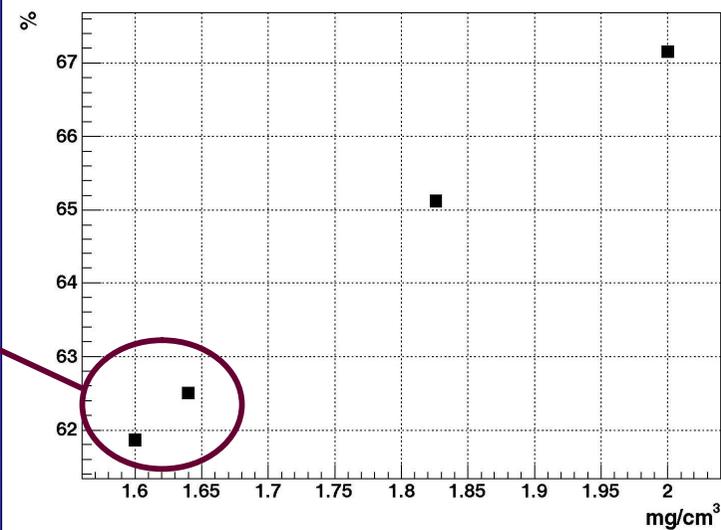


# Effect of gas density variation on the detector efficiency

Density variation



Average hardware efficiency vs density





# Solar axions

## Charge screening

**Primakoff effect**

$$\frac{d\sigma_{\alpha\leftrightarrow\gamma}}{d\Omega} = \frac{g_{\alpha\gamma\gamma}^2 Z^2 e^2}{32\pi^2} \frac{|\vec{p}_\alpha \times \vec{p}_\gamma|^2}{|\vec{p}_\alpha - \vec{p}_\gamma|^4}$$

$$V(r) = \frac{Ze}{4\pi} \cdot \frac{e^{-\kappa_s r}}{r}$$

$$\kappa_s^2 = \frac{4\pi\alpha_{em}}{T} \left( n_e + \sum_{\text{nuclei}} Z_j^2 n_j \right)$$

$$\frac{1}{\kappa_s} = \text{Debye-Huckel screening length}$$

$$x = \frac{r}{R_\odot}$$

$$\omega_p = \sqrt{\frac{4\pi\alpha_{em} n_e}{m_e}}$$

$$E^2 = k^2 + \omega_p^2$$

$$f_B = \frac{1}{e^{E/T} - 1}$$

$$\Gamma_{\gamma\rightarrow\alpha} = \frac{g_{\alpha\gamma\gamma}^2}{32\pi} \cdot T \cdot \kappa_s^2 \cdot \left[ \left( 1 + \frac{\kappa_s^2}{4E^2} \right) \cdot \ln \left( 1 + \frac{4E^2}{\kappa_s^2} \right) - 1 \right]$$

## Axion luminosity

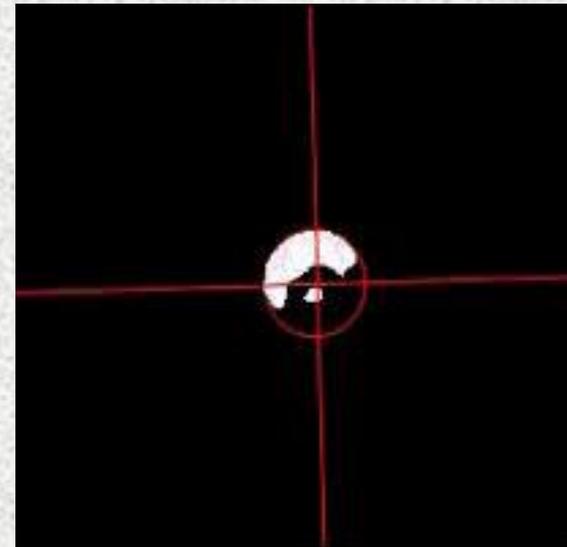
$$L_\alpha = 2R_\odot^3 \int_0^1 dx \cdot 4\pi x^2 \int_{\omega_p}^\infty dE \frac{4\pi k^2}{(2\pi)^3} \cdot \frac{dk}{dE} \cdot E \cdot f_B \cdot \Gamma_{\gamma\rightarrow\alpha}$$

## Axion flux on Earth

$$\Phi_\alpha = \frac{R_\odot^3}{\pi^2 D_\oplus^2} \int_0^1 dx \cdot x^2 \int_{\omega_p}^\infty dE \cdot k \cdot f_B \cdot \Gamma_{\gamma\rightarrow\alpha}$$



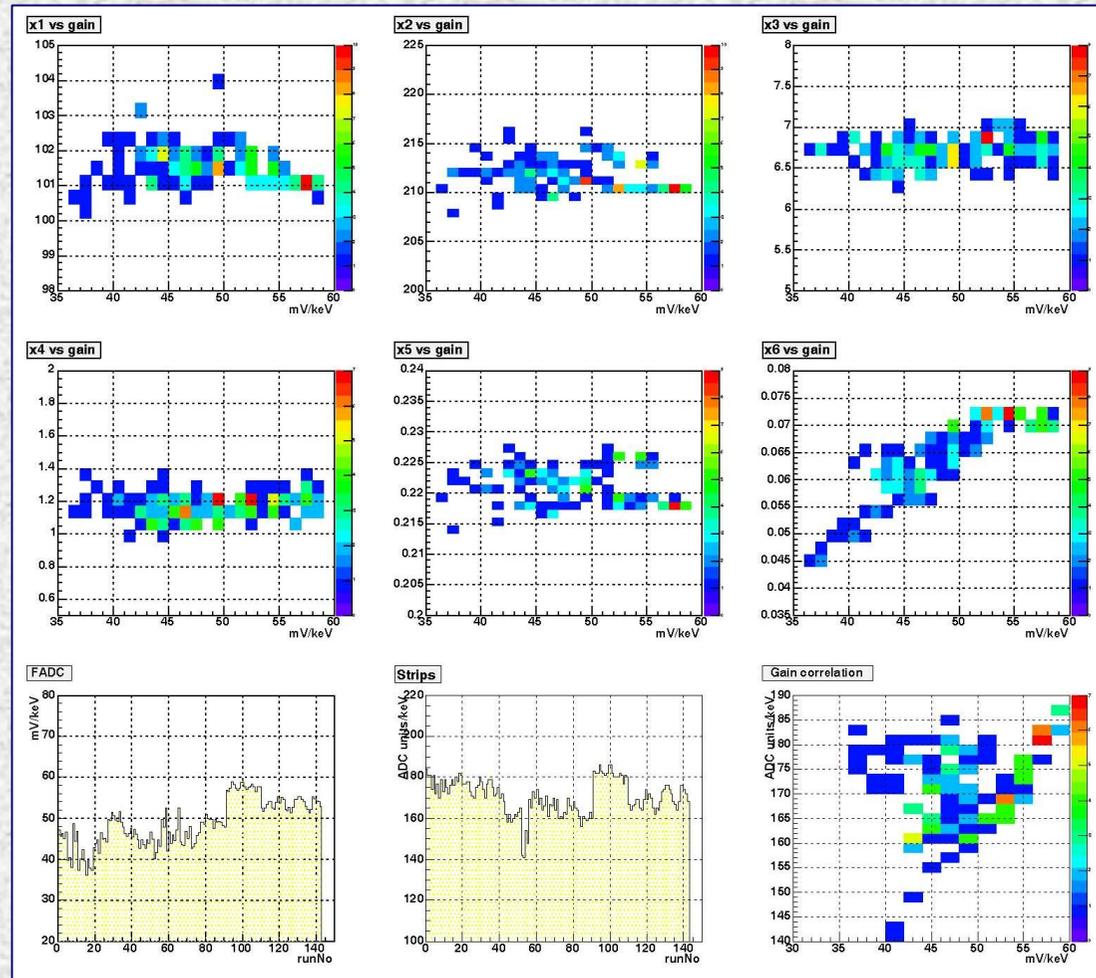
# Alignment to the Sun



March 2003



# Stability





# Stability

